

MAGNETAR GIANT FLARES AS A NEW SITE OF R-PROCESS NUCLEOSYNTHESIS

GA UK

Jakub Cehula

Charles University, Prague



Astronomical Institute, SAS & Tartu Observatory, Estonia



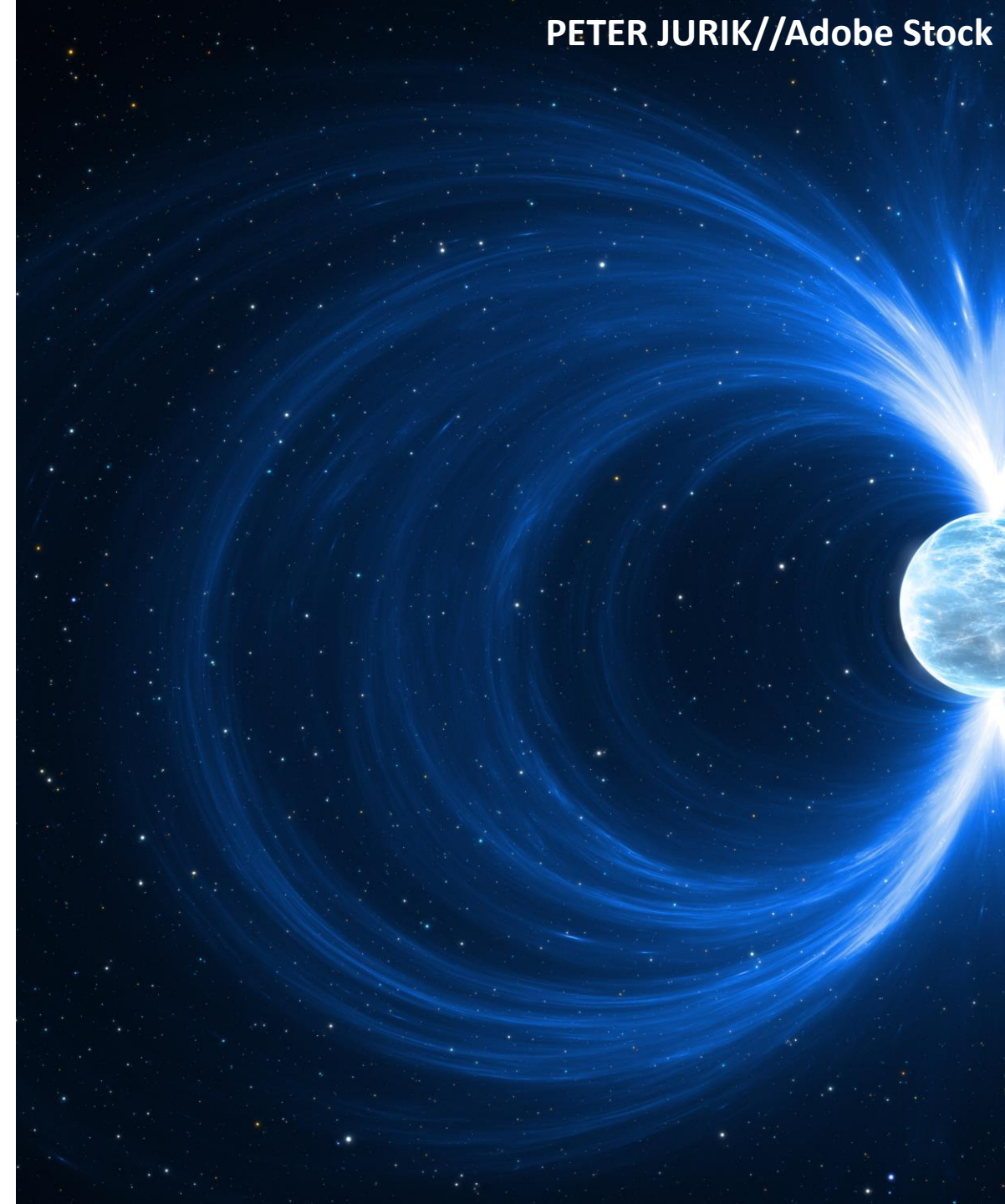
**Charles
University**

with **Anirudh Patel** (Columbia U.), Brian Metzger (Columbia U./Flatiron),
Todd Thompson (Ohio State U.), and collaborators

Tatra Astro Summit 2025, September 24, 2025

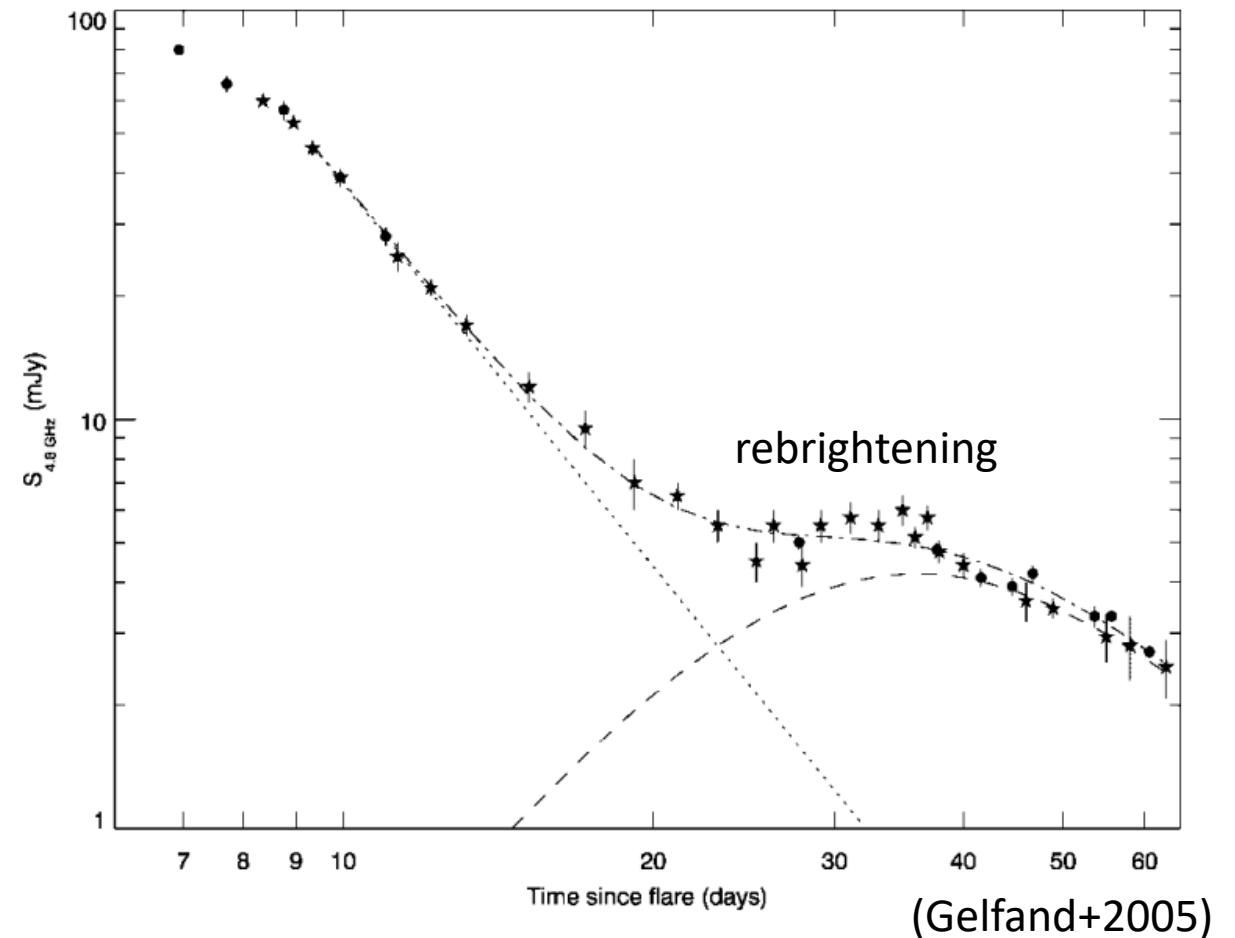
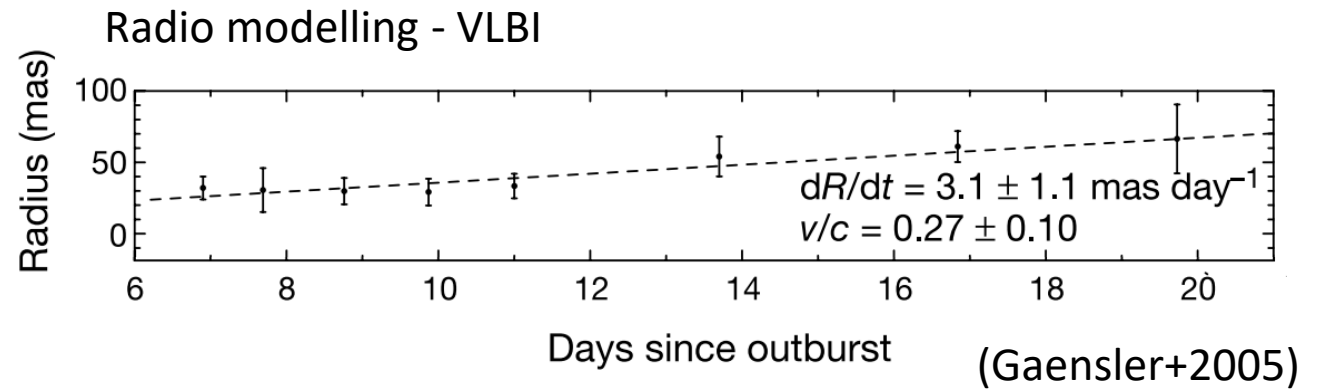
MAGNETAR GIANT FLARES

- magnetars:
 - neutron stars (NSs) with extremely strong magnetic fields ($B \gtrsim 10^{14}$ G; $B_{\odot} \sim 1$ G)
 - produced in tens of per cent of CCSNe (Beniamini+2019)
- giant flares (GFs):
 - the most powerful non-cataclysmic NS outbursts ($E \in 10^{44} - 10^{47}$ erg, $L_{\odot} \sim 4 \times 10^{33}$ erg / s)
 - galactic GFs: three in 50 years (Mazets+1979; Hurley+1999, 2005)
 - extragalactic GFs: seven candidates in 20 years (Beniamini+2024; Rodi+2025)
 - most notable GF: the **2004 GF from SGR 1806-20** (Hurley+2005)

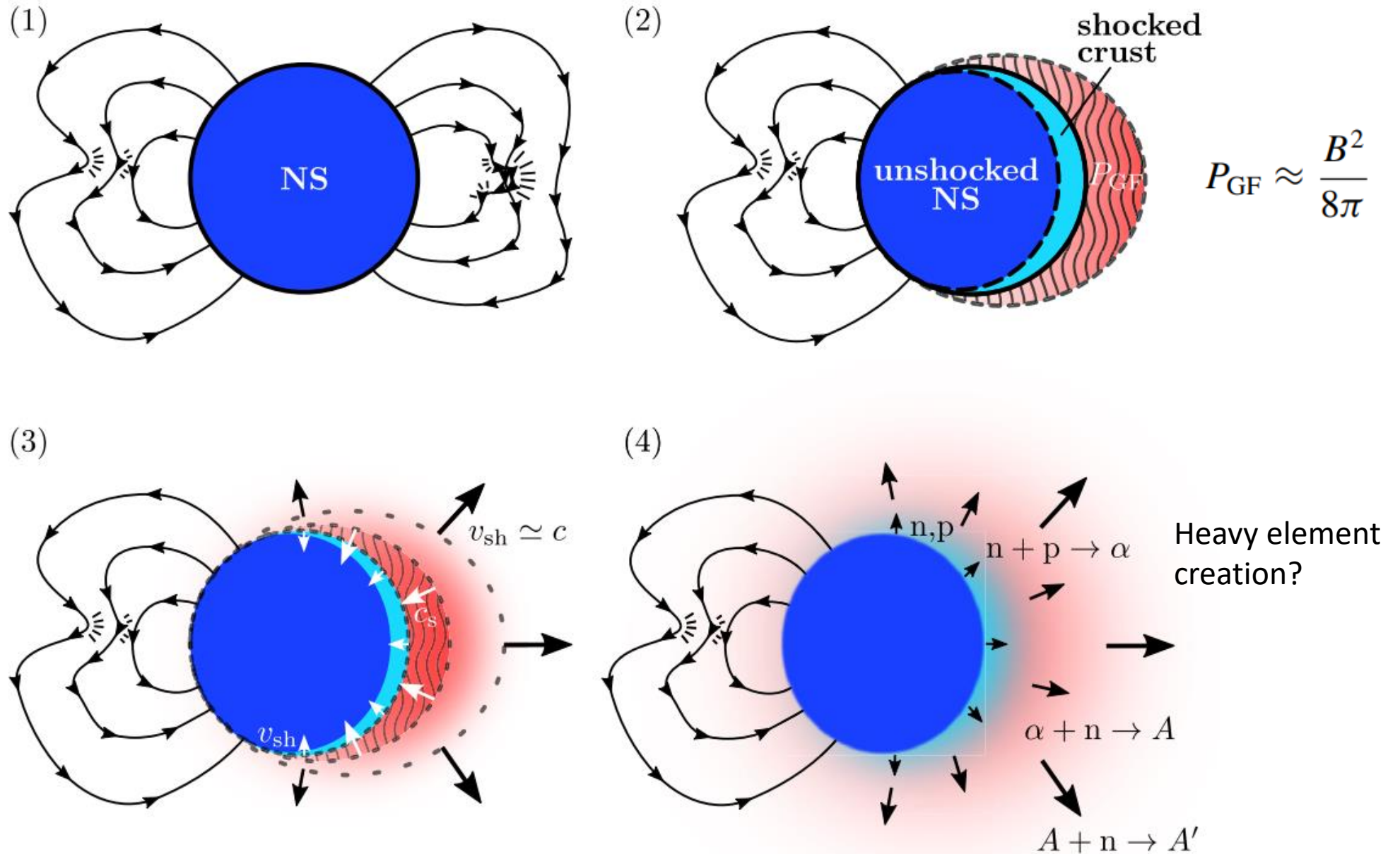


DECEMBER 2004 GF: RADIO AFTERGLOW

- GF: $\sim 2 - 4 \times 10^{46}$ erg during 0.2 – 0.5 s in soft gamma-rays and hard X-rays (Hurley+2005, Palmer+2005)
- radio afterglow (Cameron+2005; Gaensler+2005; Gelfand+2005):
 - **rebrightening** around Day 25 lasting for about a week
- modelling of the radio light curve:
 - ejecta hitting a pre-existing shell
 - mildly relativistic ($v_{ej} \sim 0.3 - 0.7c$)
ejecta of mass $M_{ej} \sim 10^{-7} M_{\odot}$
(Gelfand+2005, Granot+2006)
- **HOW** to eject $10^{-7} M_{\odot}$ from a NS?!



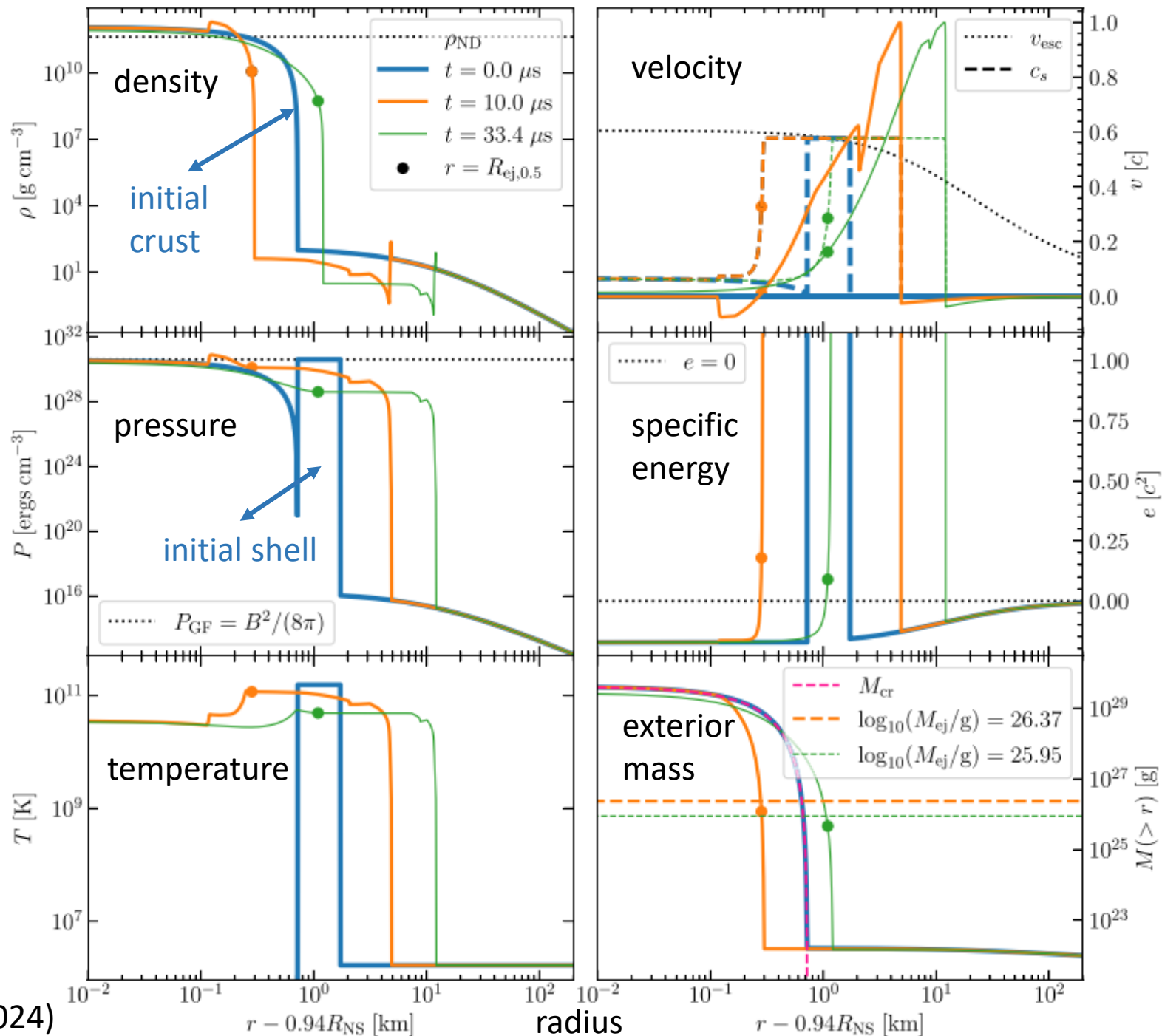
PHYSICAL PICTURE FOR BARYON EJECTION



NUMERICAL APPROACH

- 1D special-relativistic hydrodynamic simulations
- fiducial model:
 - $B = 10^{16}$ G
($P_{\text{GF}} = 10^{30.6} \text{ erg} \cdot \text{cm}^{-3}$)
 - $\Delta R_{\text{GF}} = 1$ km
- gives:
 - $M_{\text{ej}} \approx 10^{-8} - 10^{-7} M_{\odot}$
 - $\overline{v_{\text{ej}}} \approx 0.3c$
- **compatible** with the 2004 GF from SGR 1806-20 (Gelfand+2005; Granot+2006)

(PLUTO: Mignone+2007,2024)

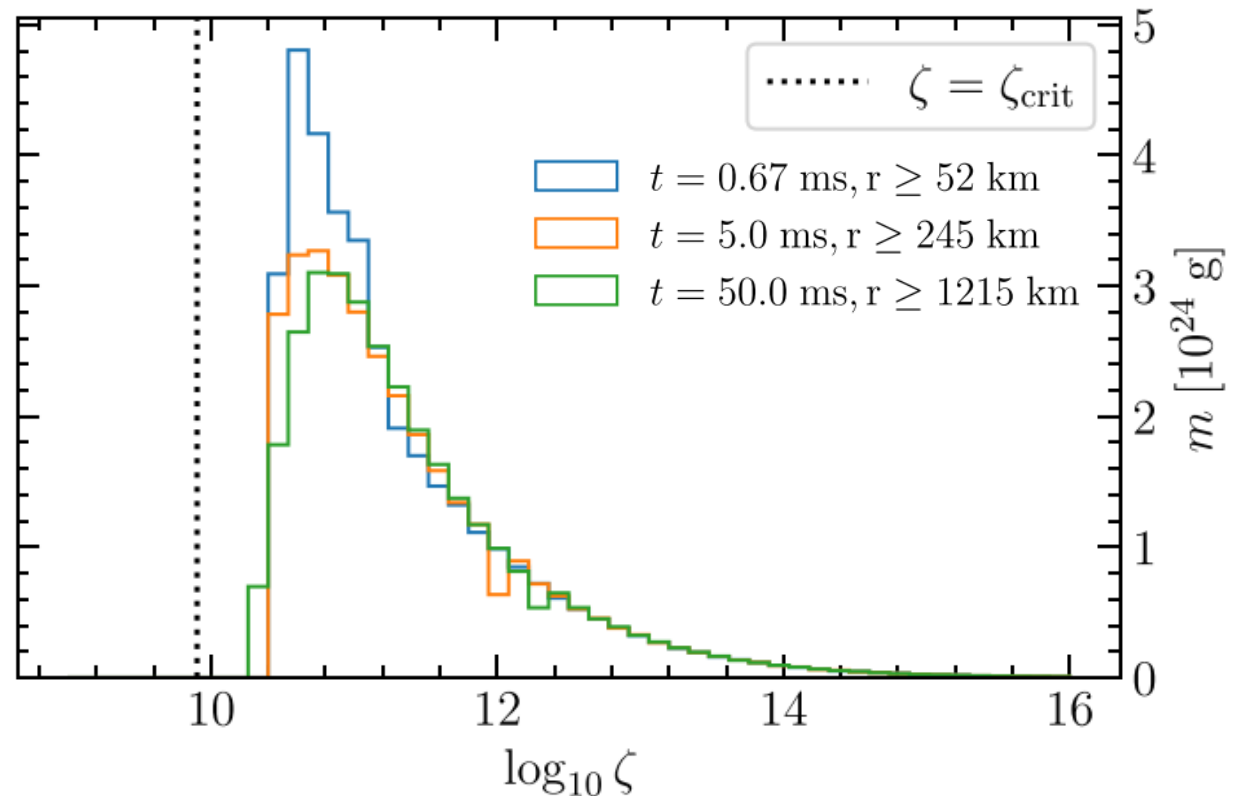


HEAVY ELEMENT CREATION

- rapid neutron-capture process (***r*-process**): a set of nuclear reactions responsible for the creation of approx. half of the atomic nuclei heavier than iron
- ***r*-process**: if $\zeta > \zeta_{\text{crit}}$, then heavy *r*-process is happening (Hoffman+1997)
- it holds:

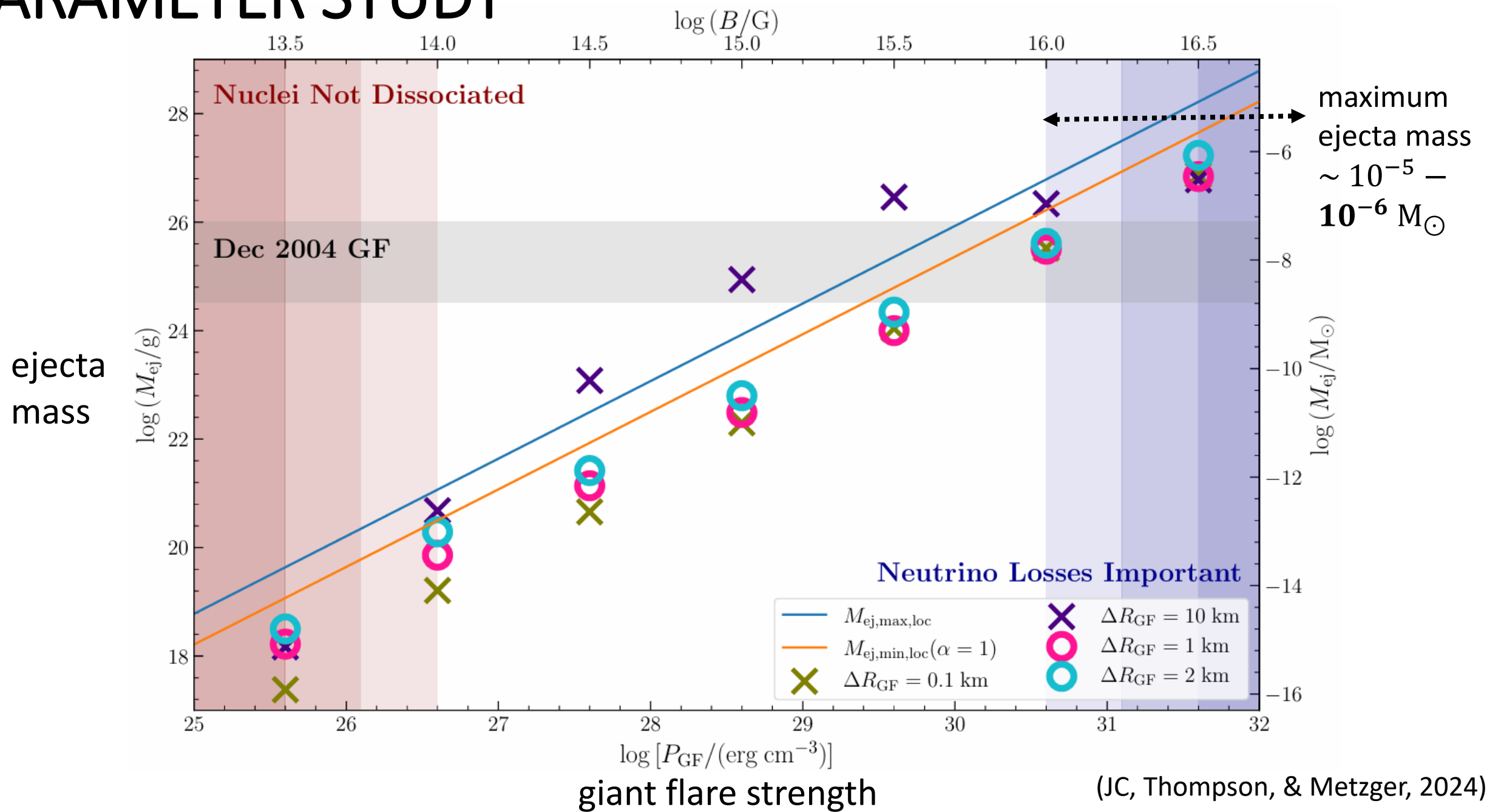
$$\zeta \equiv \frac{S_b^3}{Y_e^3 t_{\alpha,0}}$$

$$\zeta_{\text{crit}} \approx 8 \times 10^9$$



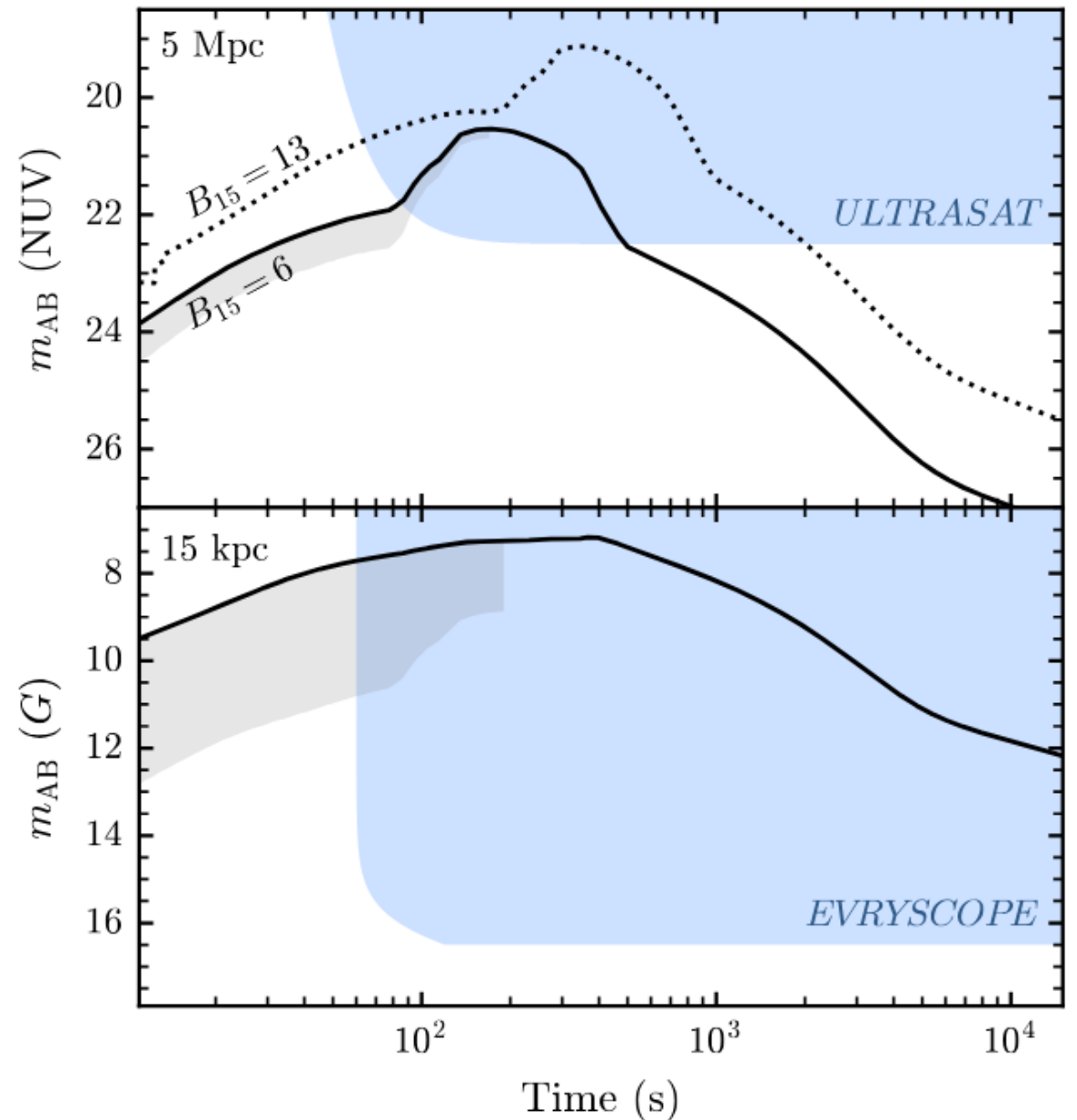
- we see that we **have heavy *r*-process!**

PARAMETER STUDY



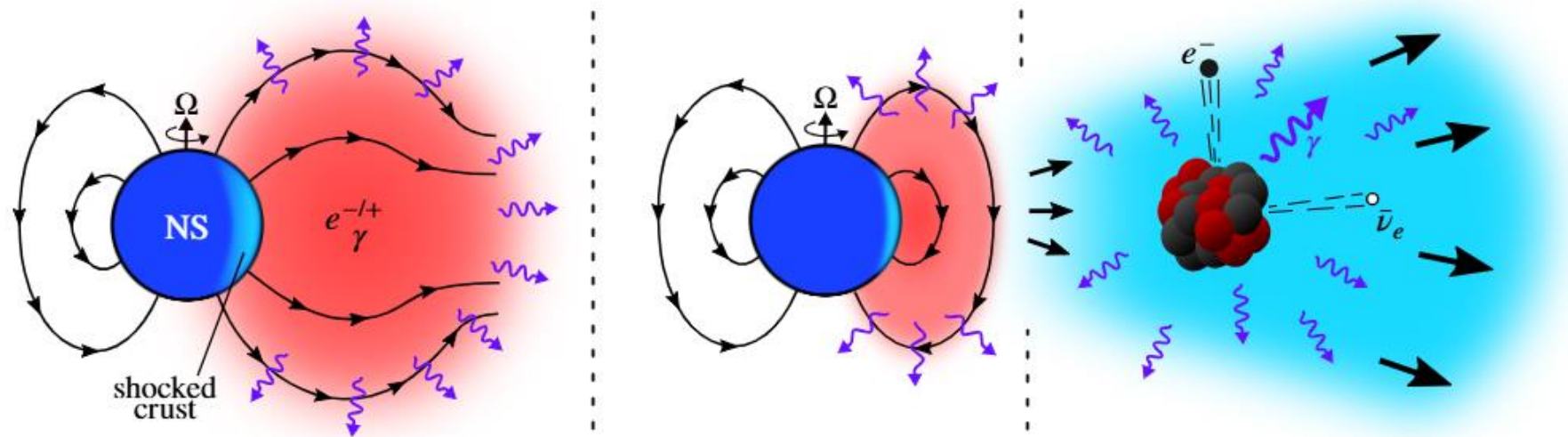
MINI-KILONOVA ($\equiv$ NOVA BREVIS) DETECTION PROSPECTS

- r -process $\Rightarrow$ radioactive heating $\Rightarrow$ kilonova-like optical/UV transient:
 - rough estimate: short $\sim$ minutes long, luminous $\sim 10^{39} \text{ erg s}^{-1}$ optical/UV transient $\equiv$ '**NOVA BREVIS**' (a brief/ short nova)
- in the Milky Way (Local Group):
 - very large instantaneous field of view: **EVRYSCOPE** (38% of the night sky with a 2-minute cadence Law+2014)
- in the nearby universe ($\lesssim 10 \text{ Mpc}$):
 - a rapidly slewing space-based optical/UV telescope: **ULTRASAT** (Sagiv+2014, launch in **2027**), UVEX (Kulkarni+2021), or QUVIK (Werner+2024)



(Patel, Metzger, Goldberg, JC, Thompson, & Renzo 2025)

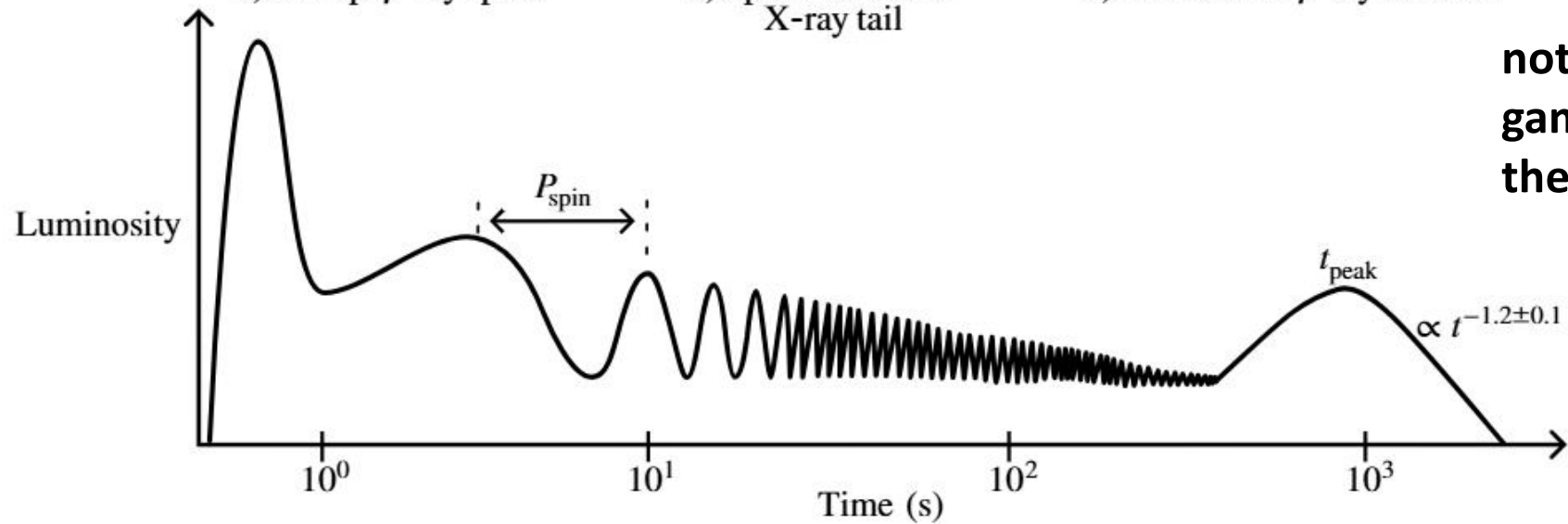
PICTURE FOR GAMMA-RAY EMISSION



a) Prompt γ -ray spike

b) Spin modulated
X-ray tail

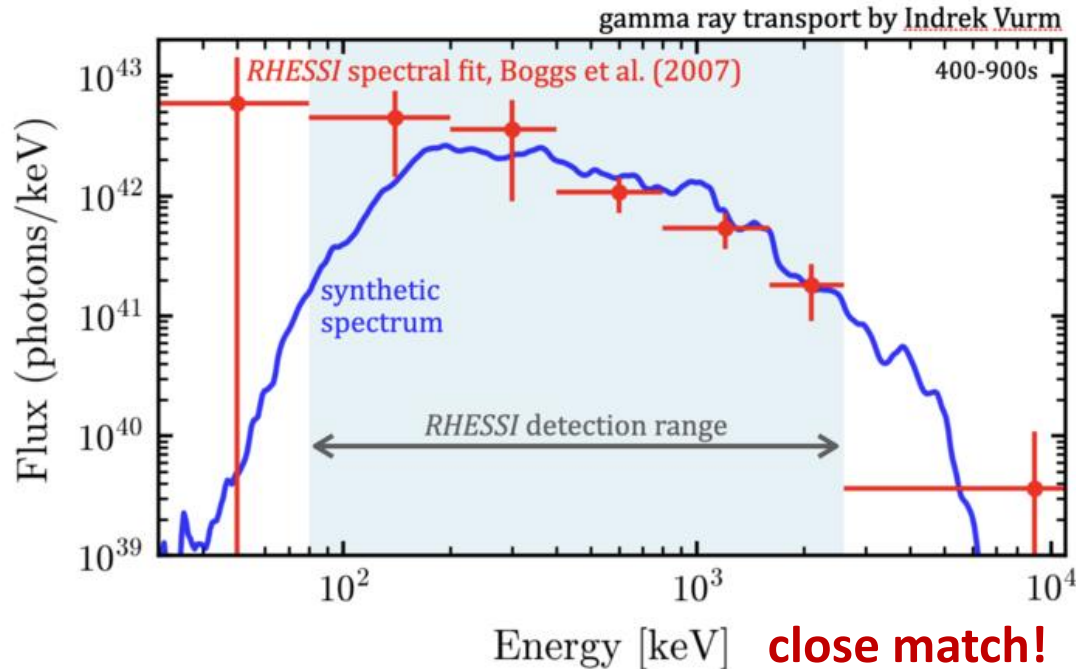
c) Radioactive γ -ray emission



not all β -decay
gamma-rays get
thermalized!

DECEMBER 2004 GF: GAMMA-RAY EMISSION

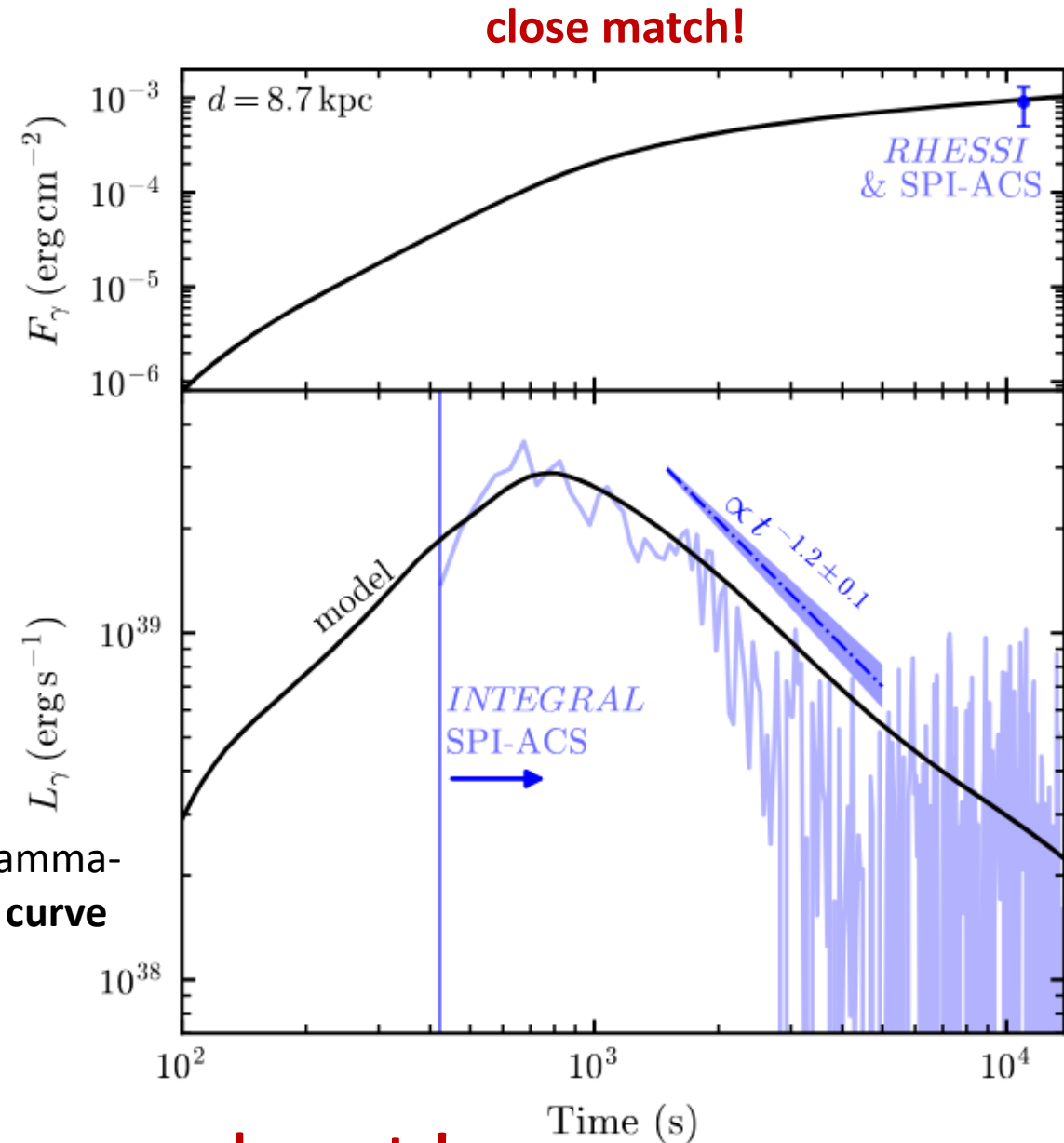
synthetic **spectra** (Doppler broadened) vs.
observations of **2004 GF** (Boggs+2007)



- conclusion: **direct evidence for $\sim 10^{-6} M_{\odot}$ of r-process elements!**

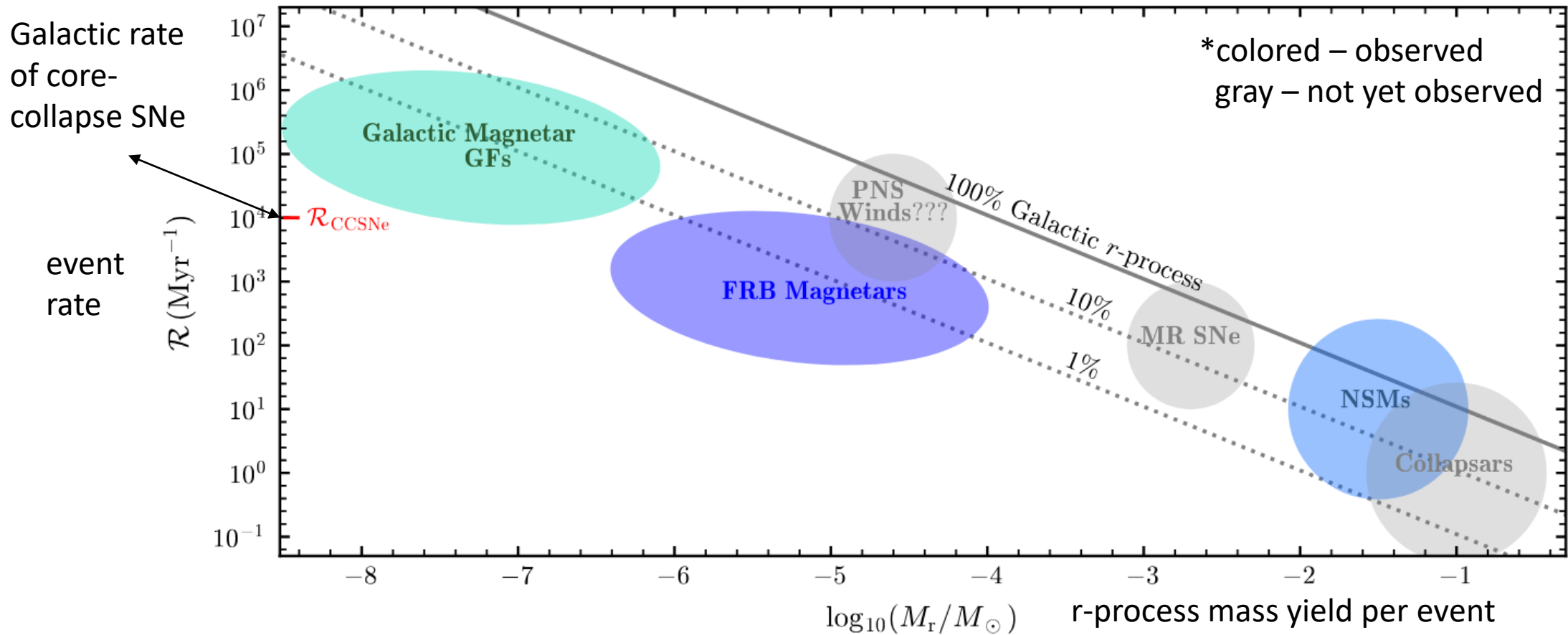
cumulative
fluence

$\sim$ MeV gamma-ray
light curve



(Patel, Metzger, JC, Burns, Goldberg, & Thompson 2025)

GALACTIC CHEMICAL EVOLUTION



- **Galactic Magnetars** could contribute **1 – 10 %** of the Galactic r -process
- FBR Magnetars are rarer but more active (e.g. Margalit+2020) => similar r -process yield
- magnetar GFs **track star formation**, some studies argue for such a channel (e.g. Lian+2023)

SUMMARY

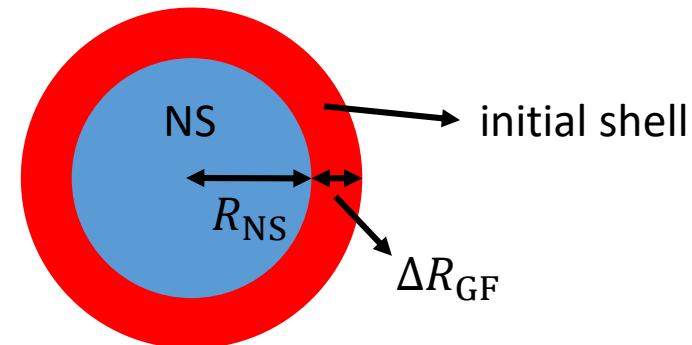
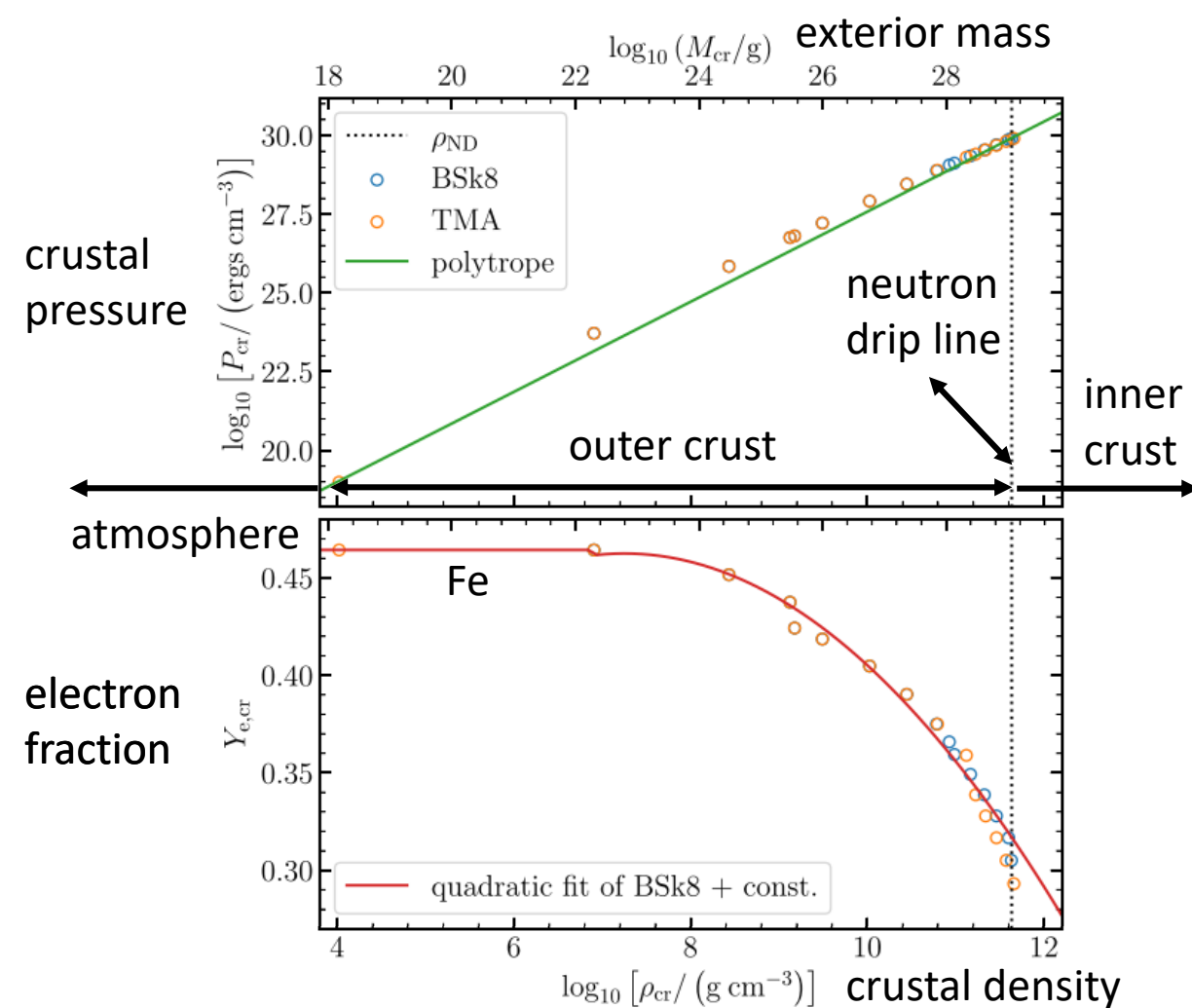
- motivation: modeling of the radio afterglow of the Dec 2004 GF from SGR 1806-20 indicates baryon ejection (Gelfand+2005, Granot+2006)
- hypothesis: GF drives a relativistic shock into the NS crust, leading to mass ejection
- methods: 1D modeling, both analytical and numerical
- results:
 - ejecta properties **align with radio** afterglow observations: mildly relativistic & $M_{\text{ej}} \sim 10^{-8} - 10^{-6} M_{\odot}$
 - r-process? **yes!**
 - radioactively powered kilonova-like transient – “**nova brevis**” (brief/short nova): $t_{\text{peak}} \sim 10 - 15$ min, $L_{\text{peak}} \sim 10^{39} - 10^{40}$ erg/s. Potentially detectable within a few Mpc (ULTRASAT/QUVIX/UVEX)
 - radioactive **gamma-ray emission**: Doppler-broadened gamma-ray lines at ~ 1 MeV, **matching an unexplained signal** following the Dec 2004 GF (**fluence, light curve, spectrum for $M_{\text{ej}} \sim 10^{-6} M_{\odot}$**)
 - **magnetar GFs** are the second **confirmed r-process site**, contributing **1-10 %** of the Galactic r-process budget. This channel tracks star formation
- future work: 2D/3D relativistic MHD simulations with neutrino cooling

(1. JC, Thompson, & Metzger, 2024; 2. Patel, Metzger, Goldberg, JC, Thompson, & Renzo 2025; 3. Patel, Metzger, JC, Burns, Goldberg, & Thompson 2025)

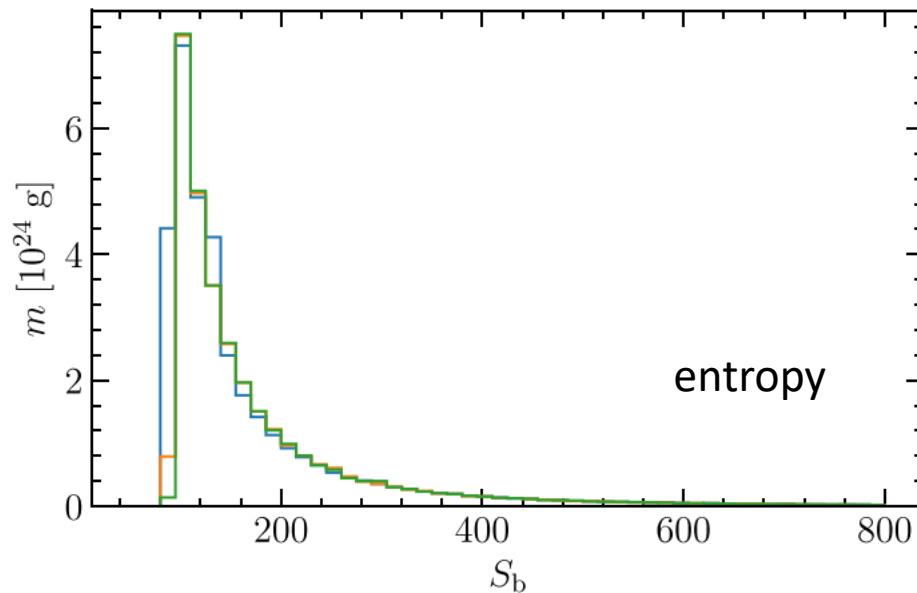
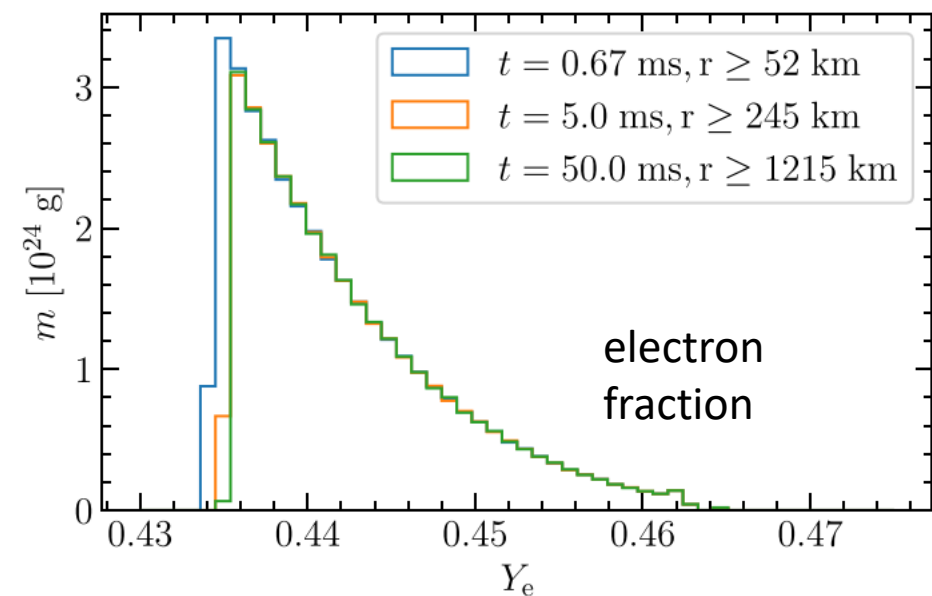
BACKUP SLIDES

NUMERICAL APPROACH

- **1D** special-relativistic hydrodynamic simulations
- crust model:
 - simple polytropic model $P_{\text{cr}} = K \rho_{\text{cr}}^{\Gamma_*}$, where $\Gamma_* = 1.43$, fits the cold (i.e. non-accreting) crust model well
- simulation set-up:
 - **initial shell** of uniform pressure $P_{\text{GF}} = B^2/8\pi$ and width $\Delta R_{\text{GF}} \leq R_{\text{NS}}$ above the NS surface
 - **no magnetic field** in the simulation
 - **ideal gas** constant- Γ EOS ($\Gamma = 4/3$ corresponding to the hot shocked gas)
 - $M_{\text{NS}} = 1.4 M_{\odot}$, $R_{\text{NS}} = 12$ km
 - reflective inner boundary: $R_{\text{in}} = 10$ km, outflow outer boundary: $R_{\text{out}} = 20000$ km



R-PROCESS? YES!

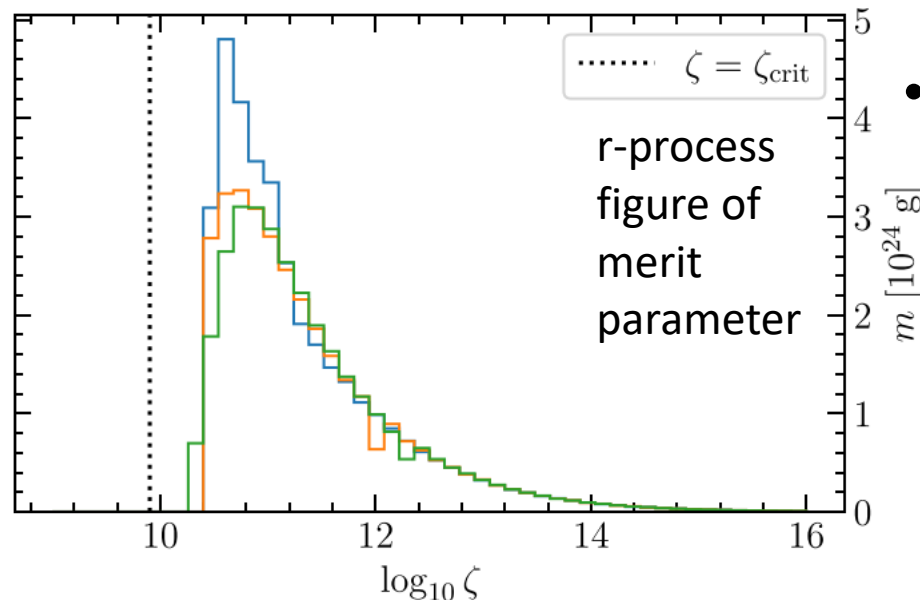
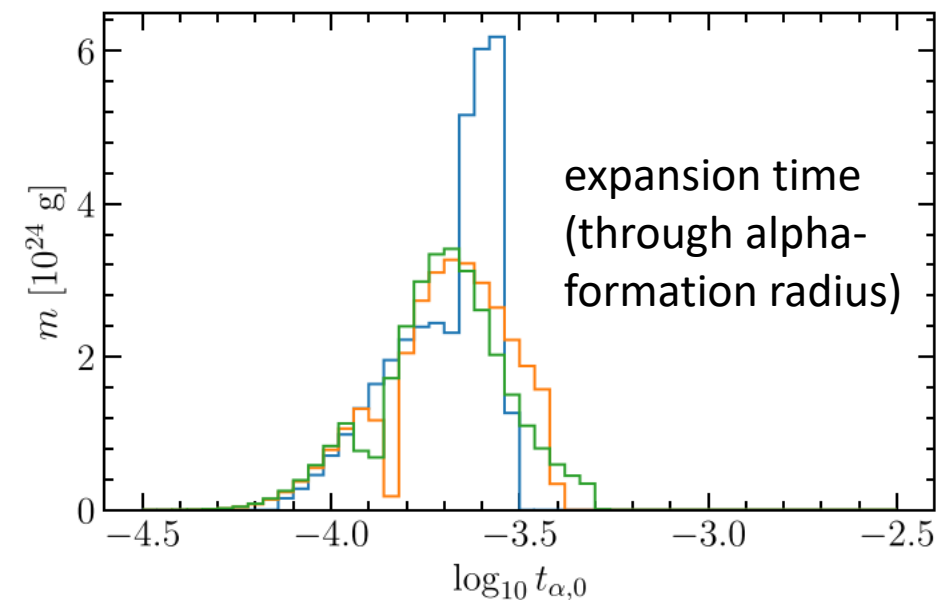


- if $\zeta > \zeta_{\text{crit}}$, then we have 3rd *r*-process peak (Hoffman+1997)

- it holds:

$$\zeta \equiv \frac{S_b^3}{Y_e^3 t_{\alpha,0}}$$

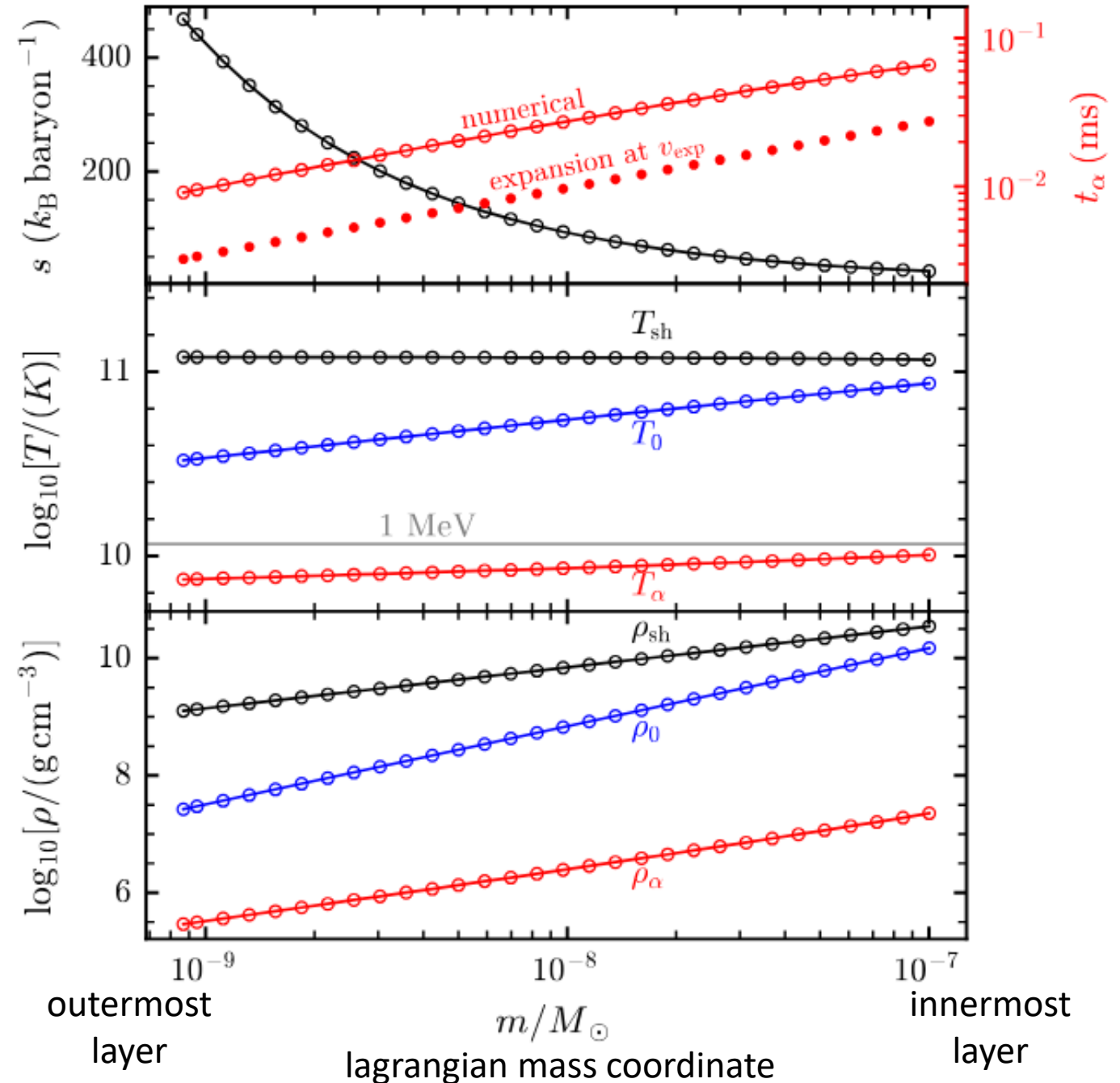
$$\zeta_{\text{crit}} \approx 8 \times 10^9$$



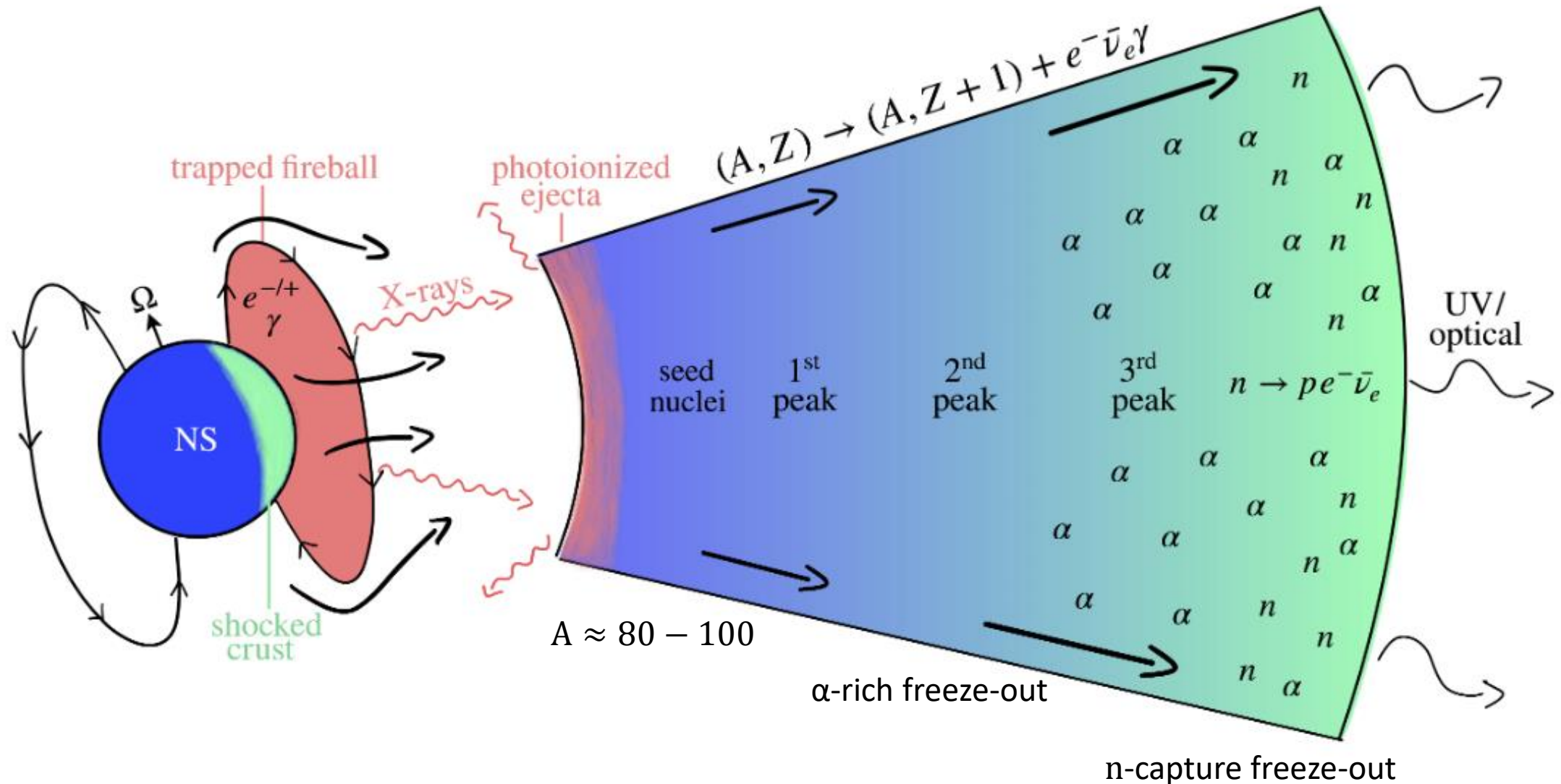
- we have 3rd *r*-process peak!

EJECTA THERMODYNAMICS

- modified Helmholtz EOS (Timmes&Swesty2000)
- constant entropy s in each layer
- fiducial model:
 - $B = 6 \times 10^{15} \text{ G}$ ($P_{\text{GF}} = 1.4 \times 10^{30} \text{ erg} \cdot \text{cm}^{-3}$) $\Rightarrow M_{\text{ej}} \sim 10^{-7} M_{\odot}$
 - $Y_e = 0.44$
 - 30 grid zones (circles)
- black – immediately post-shock
- blue – at the start of nucleosynthesis
- red – at α -formation



PICTURE FOR R-PROCESS NUCLEOSYNTHESIS



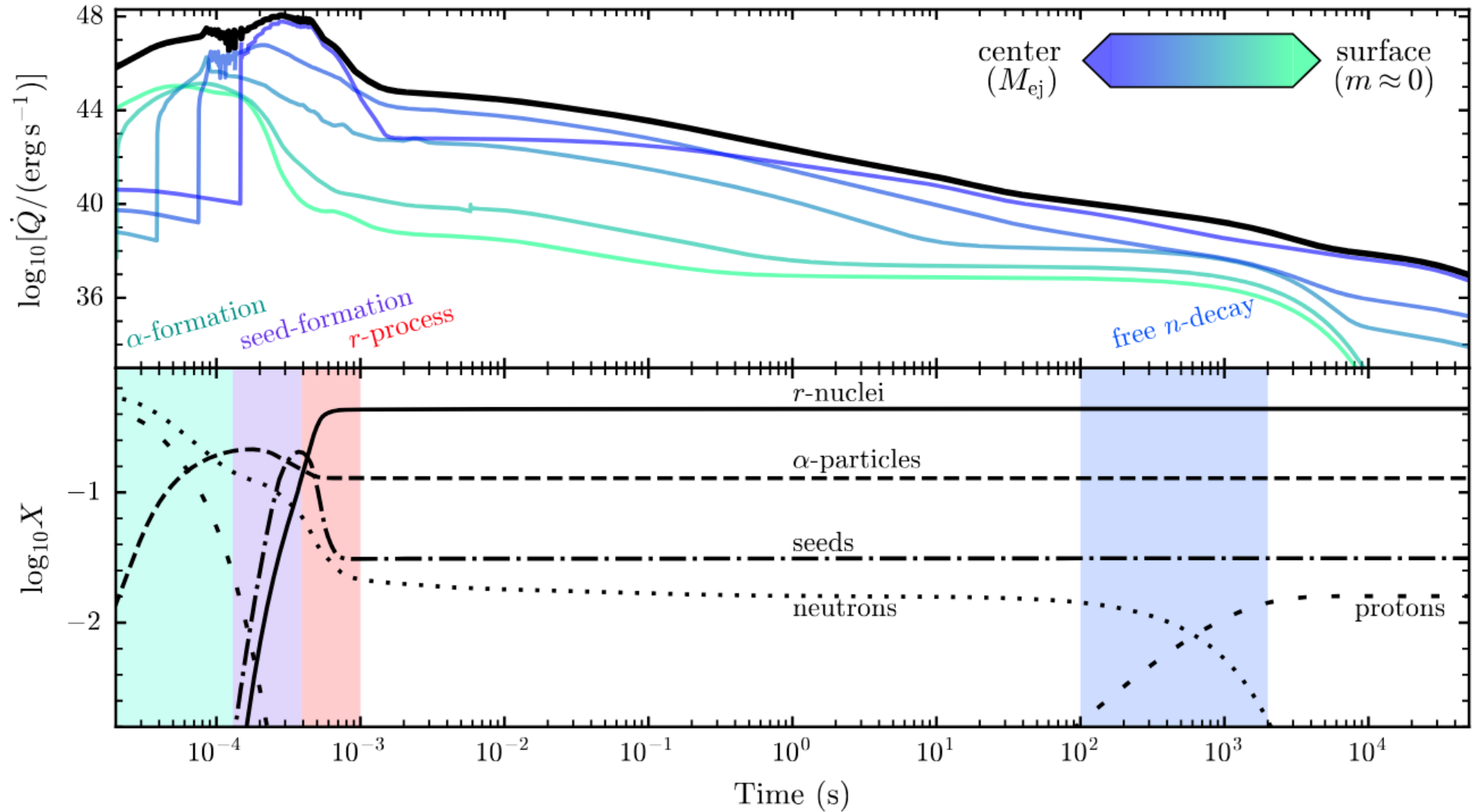
R-PROCESS NUCLEOSYNTHESIS

- initial state: $T_{\text{MeV}} \gg 1 \Rightarrow$ nucleons are free and e^-/e^+ are relativistic
- alpha-rich freeze-out mechanism:
 - 1) $T_9 \approx 10$ ($t \sim t_\alpha$): α -particle production via
 $2n + 2p \rightarrow \alpha + \gamma$,
remaining neutrons are free (n-capture freeze-out)
 - 2) $T_9 \approx 5$: carbon production via
 - 1) “normal” 3α -reactions: ${}^4\text{He}(\alpha, \gamma) {}^8\text{Be}(\alpha, \gamma) {}^{12}\text{C}$,
 - 2) neutron-aided 3α -reactions: ${}^4\text{He}(\alpha n, \gamma) {}^9\text{Be}(\alpha, \gamma) {}^{12}\text{C}$
 - 3) $T_9 \approx 2.5$: production of heavier **seed nuclei** up to the Fe-group ($A_{\text{seed}} \sim 60$) via (α, γ) reactions,
remaining α -particles are free (α -rich freeze-out)
 - 4) $T_9 \lesssim 2.5$: **r-process** proceeds to heavier nuclei via
 $(A, Z) + n \rightarrow (A + 1, Z + 1) + e^- + \bar{\nu}_e + \gamma$,
potentially up to the 3rd peak ($A \approx 190$) or beyond, depending on the
neutron to seed ratio

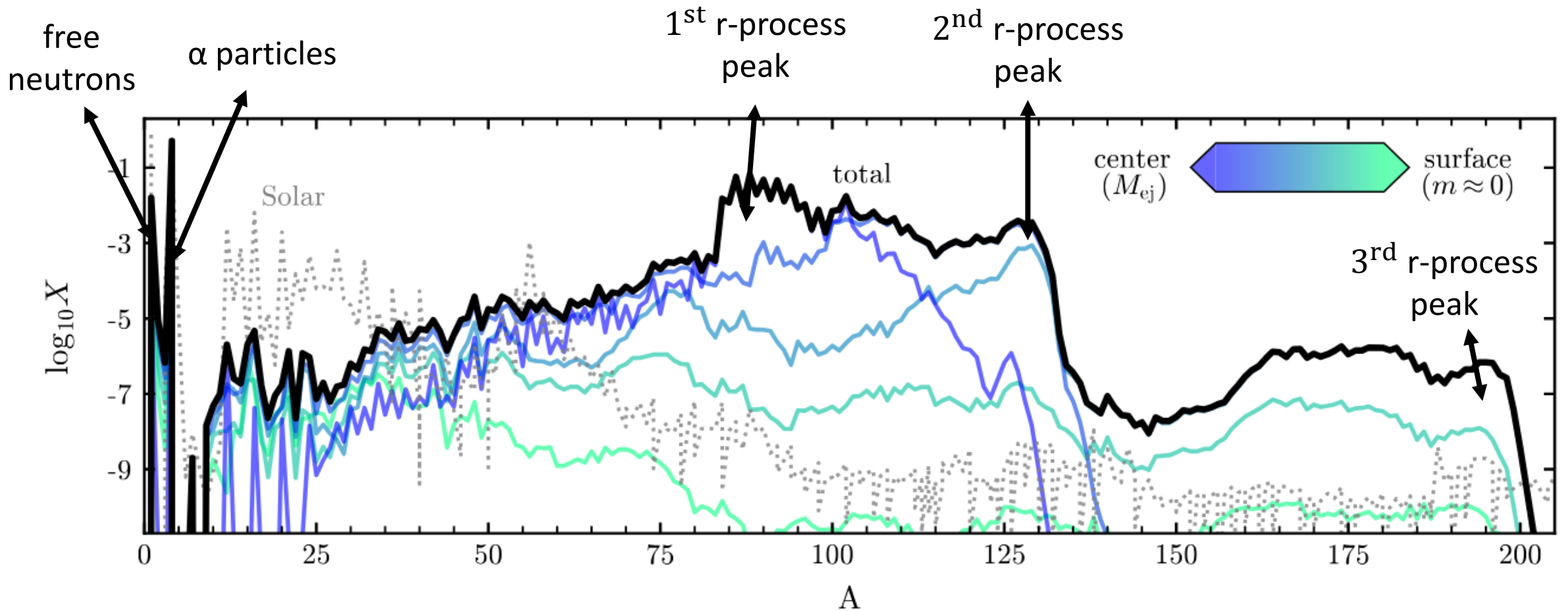
NUCLEOSYNTHESIS CALCULATIONS

total radioactive
heating rate

compositional
evolution



MASS-WEIGHTED NUCLEOSYNTHETIC YIELDS



(Patel, Metzger, Goldberg, JC, Thompson, & Renzo, accepted in ApJ)

(SkyNet: Lippuner&Roberts2015,2017; Solar System abundances: Lodders2020)

RADIOACTIVELY POWERED TRANSIENTS

- semi-analytic layer-by-layer model (Metzger2019):

➤ optical depth: $\tau(v, t) = \int_v^\infty \kappa(v', t) \rho(v', t) dv'$, with opacity given by

$$\kappa(t) = X_p(t)\kappa_p + X_\alpha\kappa_\alpha + X_{\text{seed}}\kappa_{\text{seed}} + X_{2\text{nd}}\kappa_{2\text{nd}} + X_{3\text{rd}}\kappa_{3\text{rd}},$$

➤ specific heating rate: $\dot{q}_{\text{heat}}(m, t) = f_{\text{th}}(m, t)\dot{q}(m, t)$, where the thermalization efficiency is:

$$f_{\text{th}}(t) = 0.38 \frac{\dot{q}_n(t)}{\dot{q}(t)} + \frac{\dot{q}_r(t)}{\dot{q}(t)} \left[0.25 + 0.4 \{1 - \exp(-\tau(t)\kappa_\gamma/\kappa(t))\} \right],$$

$\nwarrow$ $\nearrow$ $\nwarrow$ $\nearrow$ $\nwarrow$ $\nearrow$
 β -decay e^- neutron heating r-process heating β -decay e^- β -decay gamma-rays

➤ diffusion surface: $\tau(r) = \frac{c}{v} \Rightarrow M_{\text{diff}}$, and photosphere: $\tau(r) = \frac{2}{3} \Rightarrow M_{\text{ph}}$

➤ Arnett's law (Arnett1980):

1) optically thick (“**photospheric**”) emission: $L_{\text{ph}} = \int_{M_{\text{ph}}}^{M_{\text{diff}}} \dot{q}_{\text{heat}} dm$

2) optically thin (“**nebular**”) emission: $L_{\text{neb}} = \int_0^{M_{\text{ph}}} \dot{q}_{\text{heat}} dm$

MINI-KILONOVA

("nova brevis")

radioactive heating
(with thermalization
correction)

outermost layer

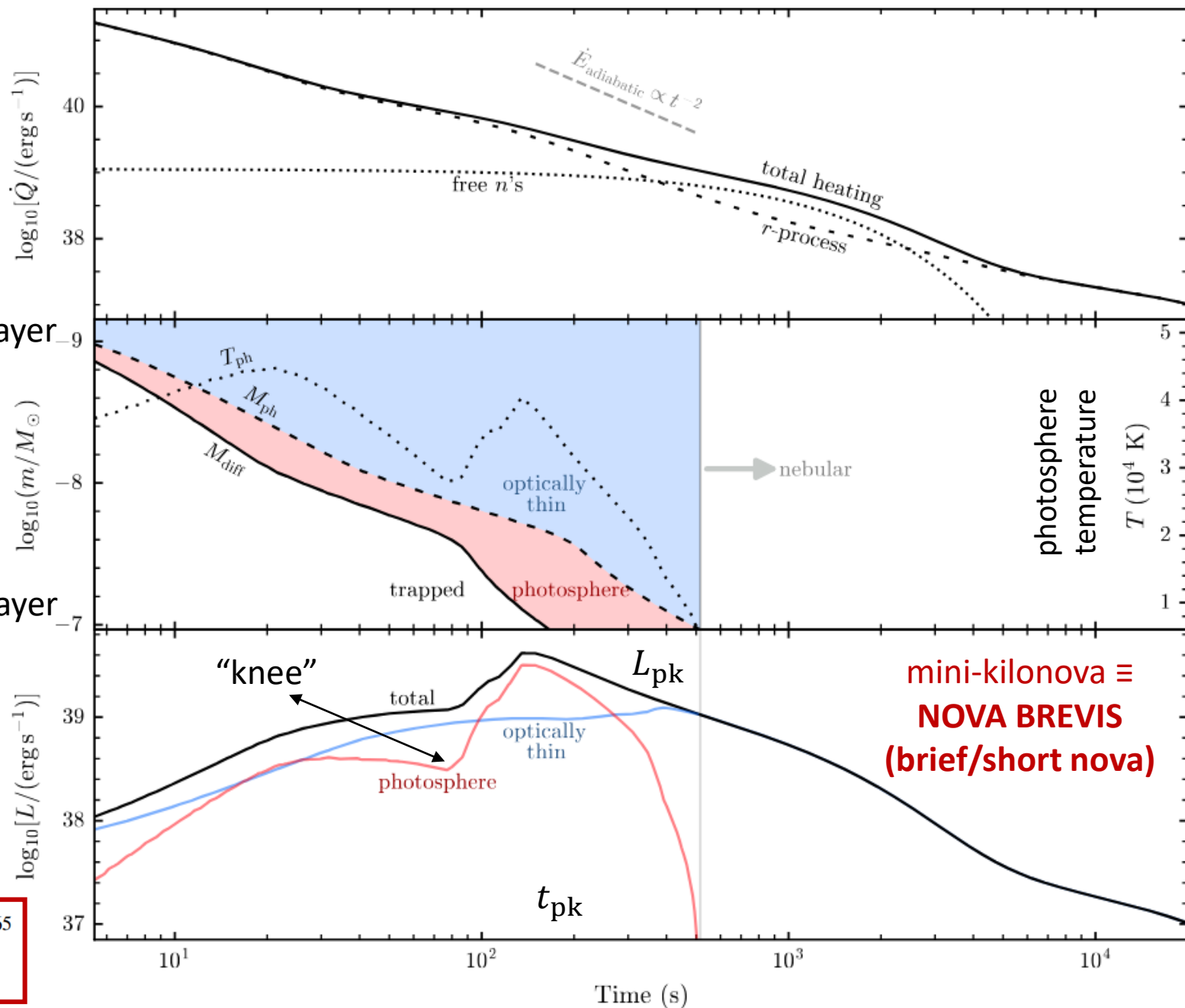
lagrangian mass
coordinate

innermost layer

luminosity

$$t_{\text{pk}} \approx 300 \text{s} \left(\frac{M_{\text{ej}}}{10^{26} \text{g}} \right)^{1/2} \left(\frac{v_{\text{ej}}}{0.3c} \right)^{-1/2} \left(\frac{\kappa}{3 \text{ cm}^2 \text{ g}^{-1}} \right)^{1/2},$$

$$L_{\text{pk}} \approx 10^{39} \text{ erg s}^{-1} \left(\frac{M_{\text{ej}}}{10^{26} \text{g}} \right)^{0.35} \left(\frac{v_{\text{ej}}}{0.3c} \right)^{0.65} \left(\frac{\kappa}{3 \text{ cm}^2 \text{ g}^{-1}} \right)^{-0.65}$$



NOVA BREVIS SPECTRA

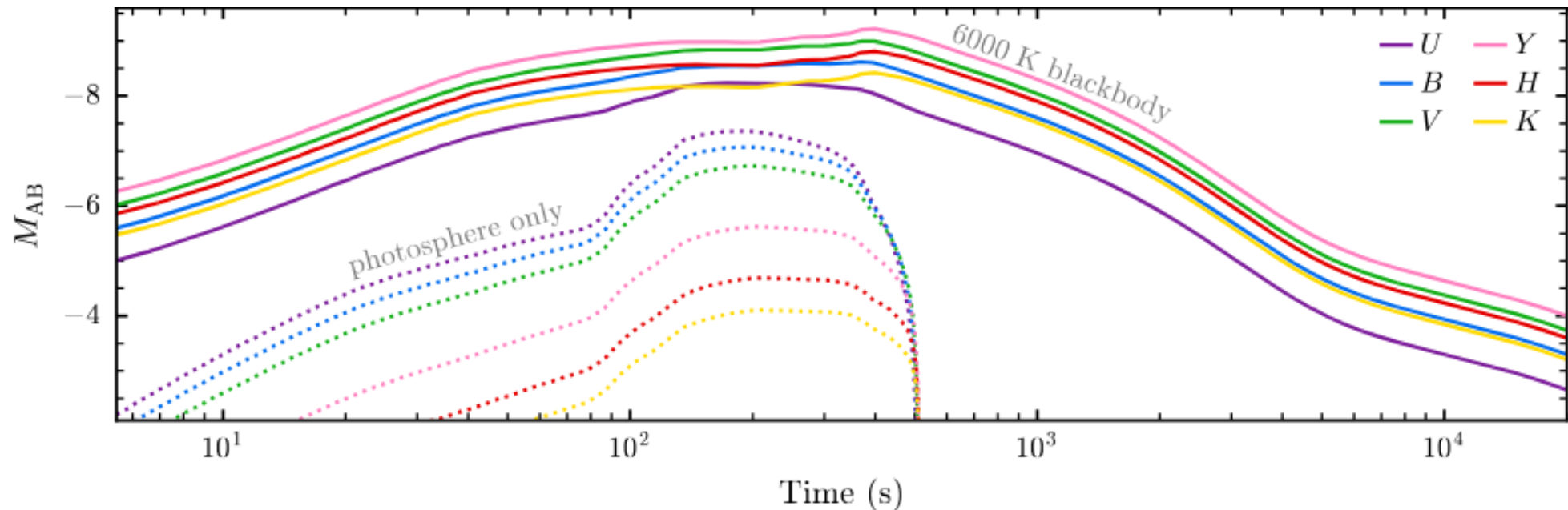
- two blackbody radiators:

➤ optically thick (“photospheric”) emission: $T_{\text{ph}} = \left(\frac{L_{\text{ph}}}{4\pi\sigma R_{\text{ph}}^2} \right)^{1/4},$

➤ optically thin (“nebular”) emission: $T_{\text{eff}} = 6000 \text{ K}$

- total flux: $F_{\nu} = F_{\nu, \text{ph}} + F_{\nu, \text{neb}} = \frac{1}{4d^2\sigma_{\text{SB}}} \left[\frac{L_{\text{ph}}}{T_{\text{ph}}^4} B_{\nu}(T_{\text{ph}}) + \frac{L_{\text{neb}}}{T_{\text{eff}}^4} B_{\nu}(T_{\text{eff}}) \right]$

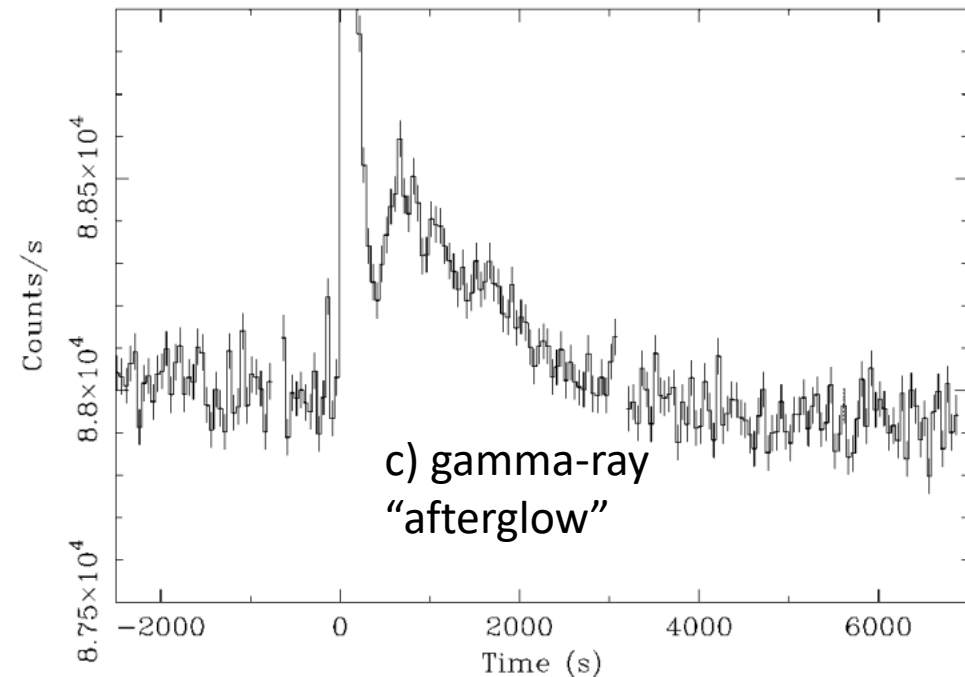
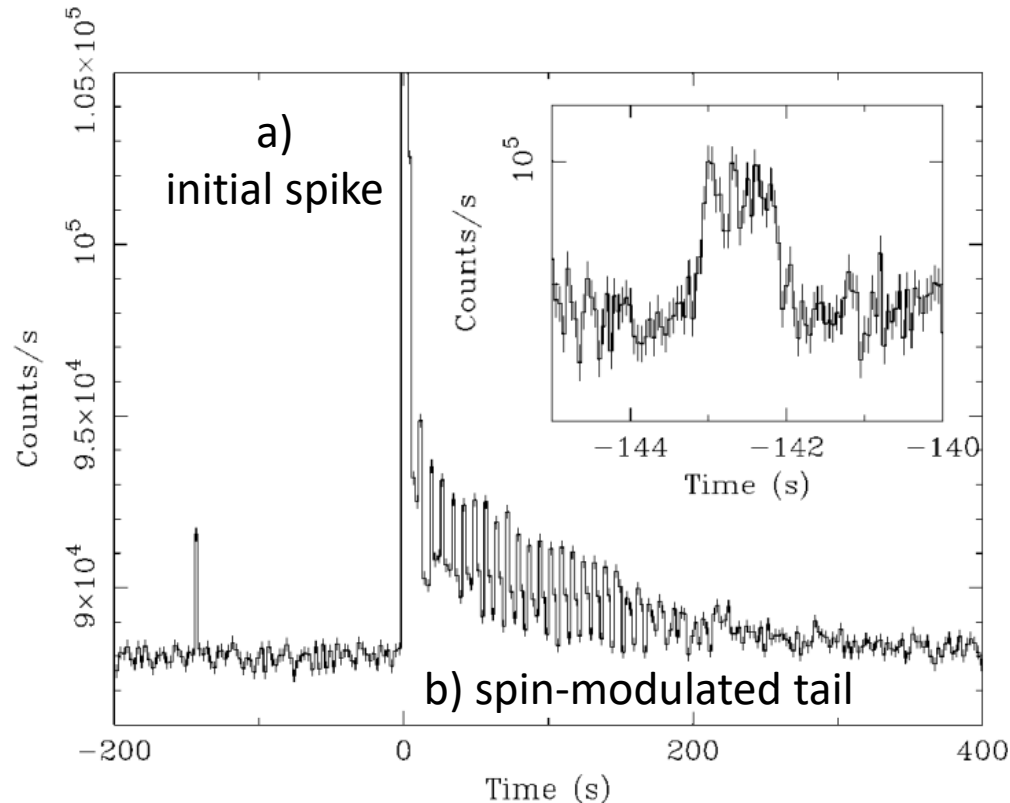
absolute AB
magnitudes in UV,
optical, and
infrared bands



DEC 2004 GF FROM SGR 1806-20

- a) $\sim 2 - 4 \times 10^{46}$ erg in gamma rays during the 0.2 – 0.5 s long initial spike
- b) $\sim 10^{44}$ erg in decaying minutes-long tail modulated at the NS period
- c) $\sim 10^{43}$ erg in hour-long gamma-ray “afterglow” that peaked after ~ 600 – 800 s

(Hurley+2005; Palmer+2005; Mereghetti+2005; Boggs+2007; Frederiks+2007)

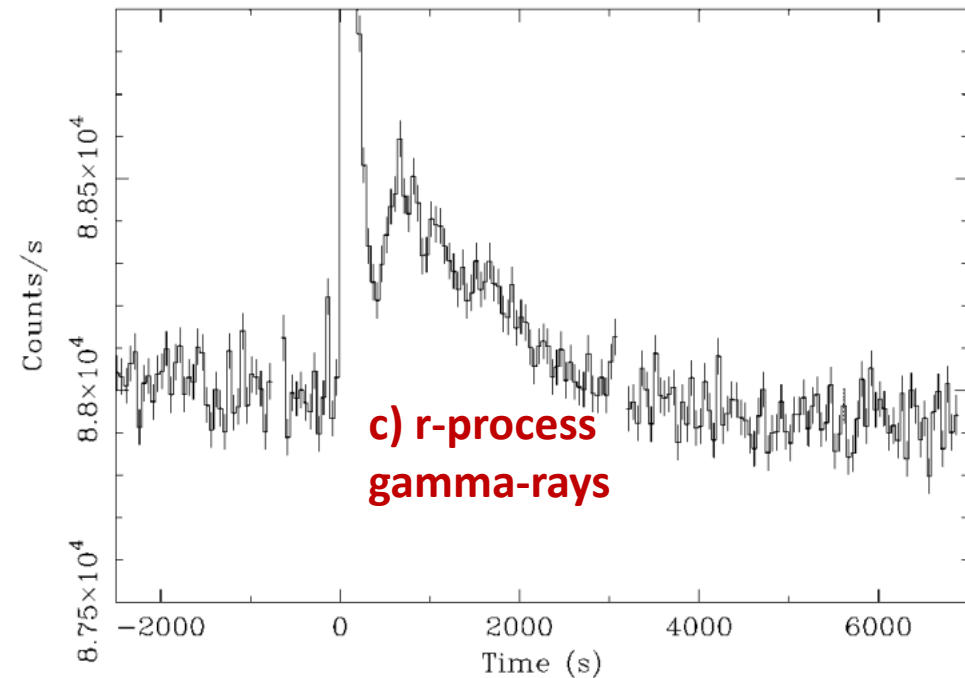
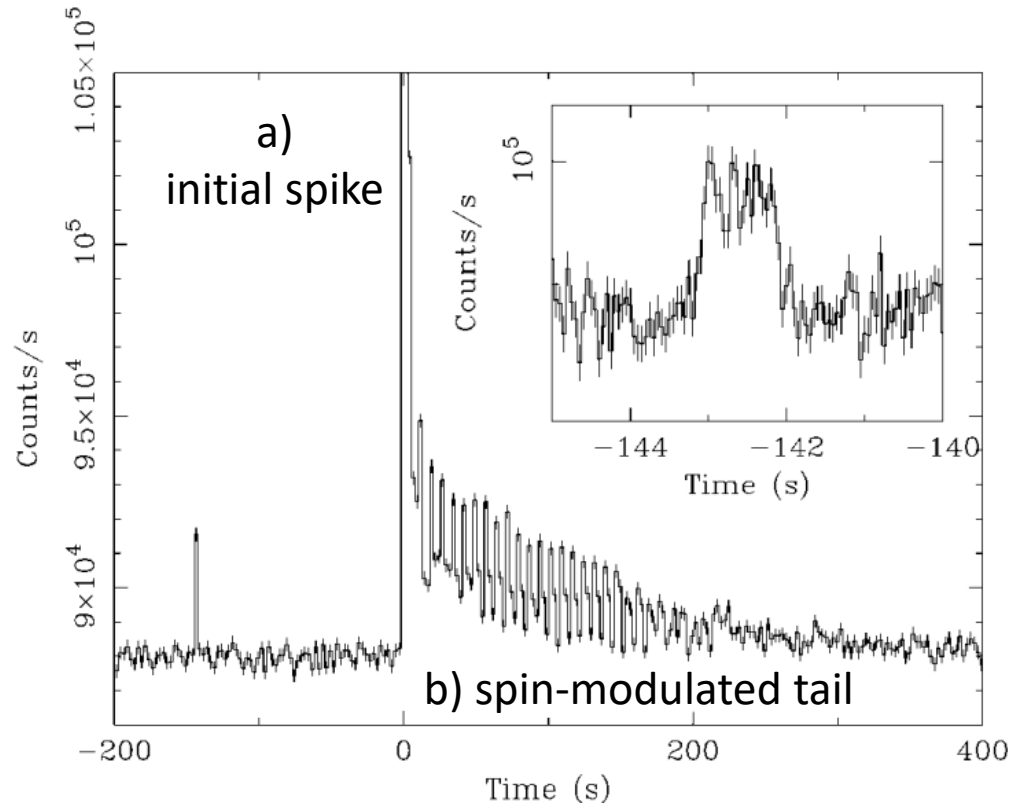


(Mereghetti+2005)

DEC 2004 GF FROM SGR 1806-20 **REVISITED**

- a) $\sim 2 - 4 \times 10^{46}$ erg in gamma rays during the 0.2 – 0.5 s long initial spike
- b) $\sim 10^{44}$ erg in decaying minutes-long tail modulated at the NS period
- **c) $\sim 10^{43}$ erg in hour-long delayed gamma-ray emission that peaked after $\sim 600-800$ s**

(Hurley+2005; Palmer+2005; Mereghetti+2005; Boggs+2007; Frederiks+2007; **Patel+2025**)



(Mereghetti+2005)

GAMMA-RAY EMISSION

- the same layer-by-layer approach
- approximate gamma-ray thermalization treatment of Hotokezaka+2016
- assuming constant opacity $\kappa_\gamma = 0.1 \text{ cm}^2 \cdot \text{g}^{-1}$
- considering all nuclear species that contribute at least 1% between 10 s and 14 000 s
- effective decay mode: $i \rightarrow f + e^- + \bar{\nu}_e + \gamma$
- weak reaction rates from JINA REACLIB database (Cyburt+2010)

- fiducial model:

$$\triangleright M_{\text{ej}} = 1.2 \times 10^{-6} M_\odot \quad (M_r = 7 \times 10^{-7} M_\odot)$$

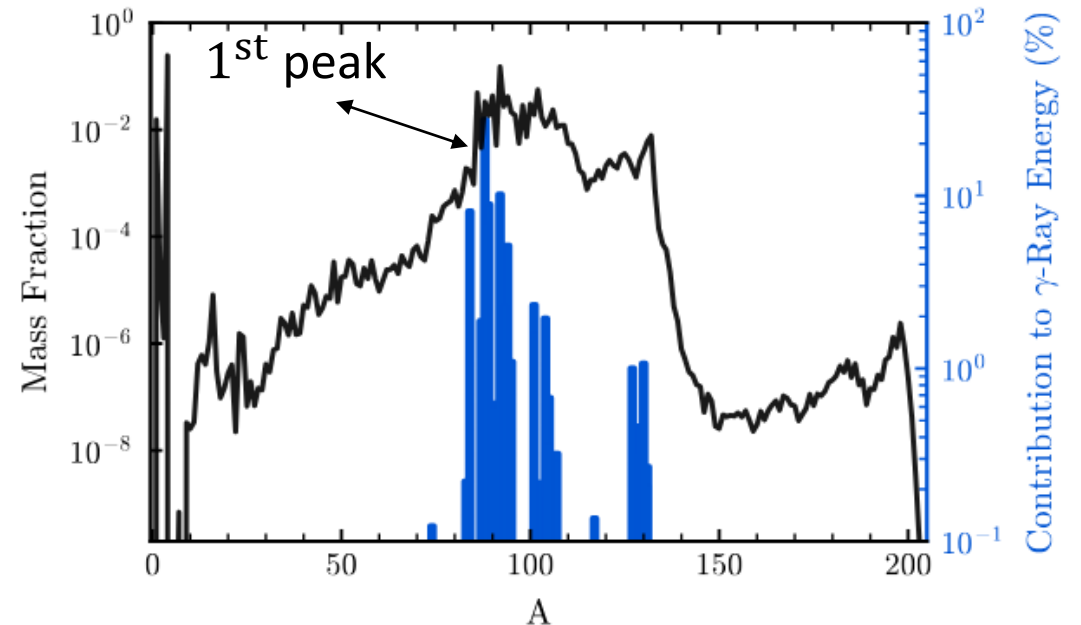
$$\triangleright Y_e = 0.40$$

$$\triangleright \bar{v} = 0.15c$$

$$\triangleright \beta = 6$$

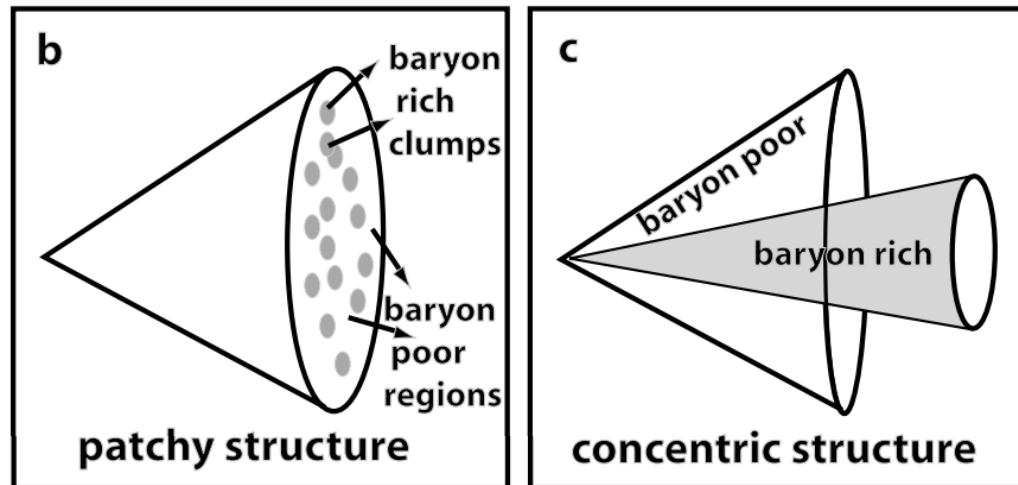
$$\triangleright f_\Omega = 3/4$$

$$\rho(r, t) = \frac{\beta - 3}{4\pi f_\Omega} \frac{M_{\text{ej}}}{(\bar{v}t)^3} \left(\frac{v}{\bar{v}}\right)^{-\beta}, v > \bar{v},$$

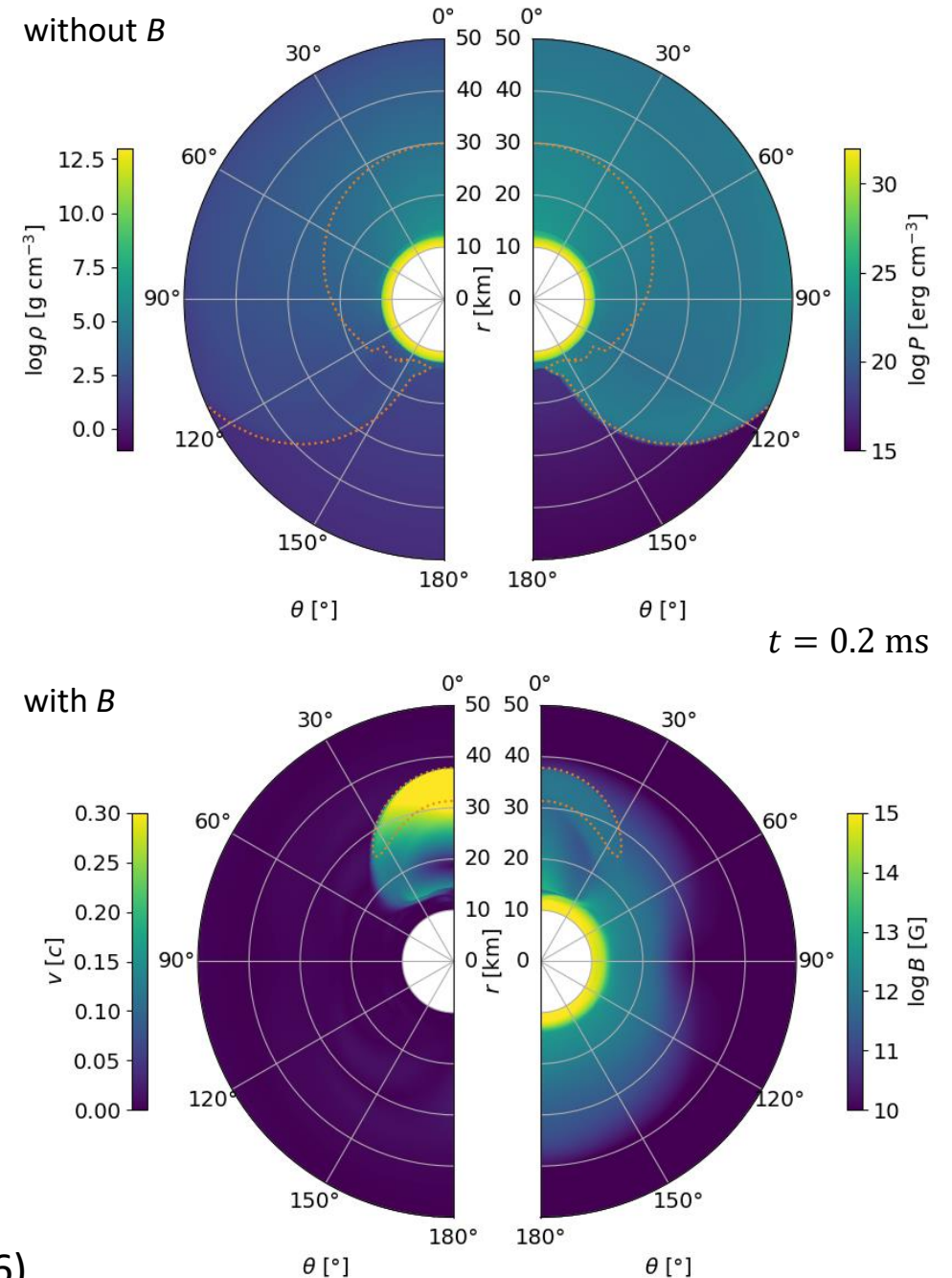


CURRENT WORK

- 2D special-relativistic MHD simulations
- motivation: magnetic field confines ejecta, preventing full **surface obscuration**, as ruled out by observations (Gelfand+2005, Granot+2006)
- preliminary result: the ejecta is more confined

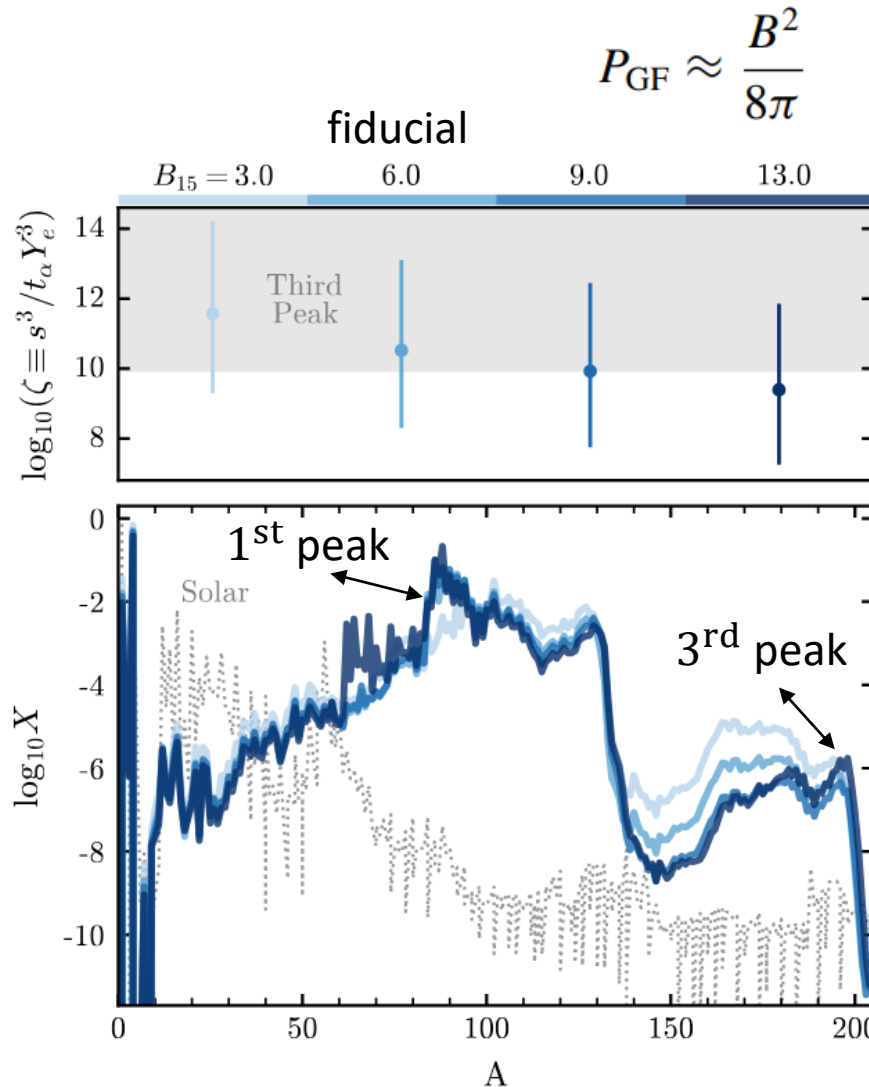


(Granot+2006)



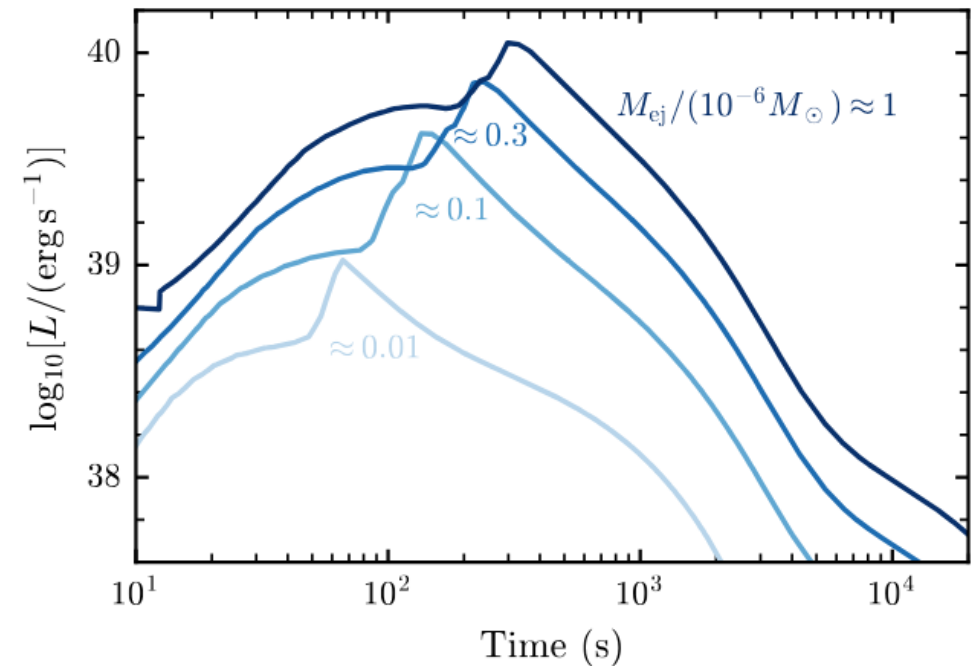
IMPACT OF VARYING THE FLARE STRENGTH

r-process
figure of
merit
parameter



mass-
weighted
nucleo-
synthesis
yields

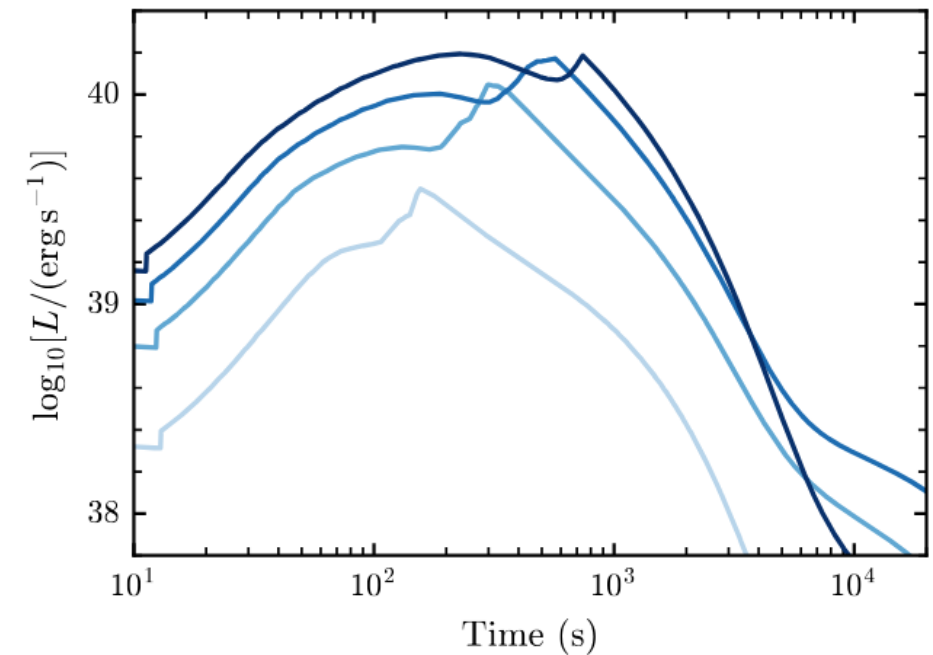
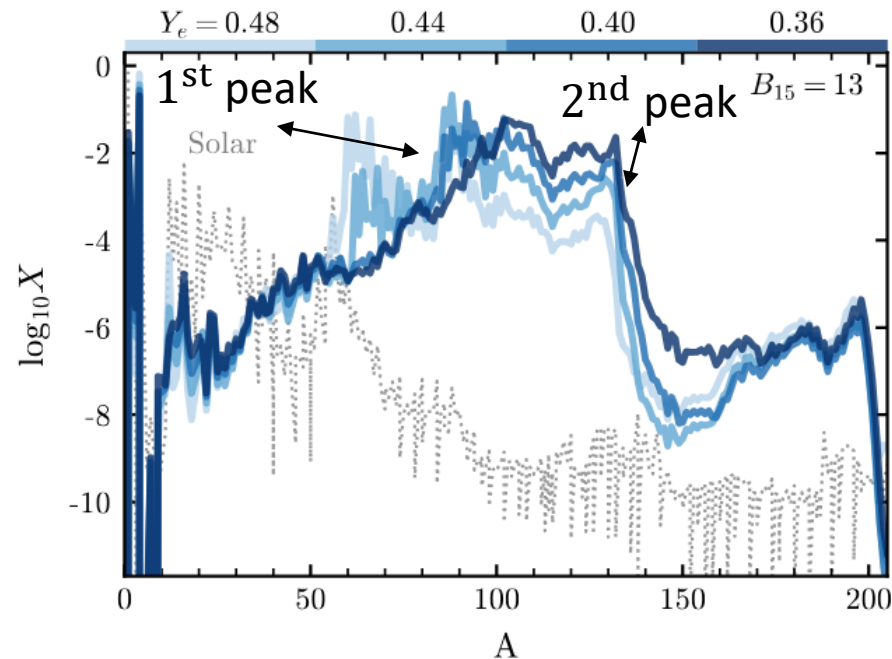
- $M_{\text{ej}}/M_{\odot} \in (10^{-8} - 10^{-6})$ for $B_{15} \in (3, 13)$
- X_r increases with flare strength
- stronger flares produce brighter and longer duration transients



IMPACT OF VARYING THE ELECTRON FRACTION

- for $B_{15} = 13$ ($M_{\text{ej}} \sim 10^{-6} M_{\odot}$) we explore $Y_e = 0.36, 0.40, 0.44, 0.48$
- X_r decreases with Y_e
- lighter elements with increasing Y_e
- higher X_r increases luminosity, different composition affects opacity $\Rightarrow t_{\text{pk}}$

mass-
weighted
nucleo-
synthesis
yields



IMPACT OF VARYING THE MODEL PARAMETERS

- larger Y_e :
 - lower $X_r \Rightarrow$ reducing the gamma-ray emission
- larger $\bar{v}$:
 - ejecta becomes optically-thin earlier when the energy generation rate is larger $\Rightarrow$ earlier and brighter peak emission
- larger β :
 - less mass in the outer ejecta layers $\Rightarrow$ steeper light-curve rise
- larger M_{ej} :
 - increases the peak timescale and gamma-ray luminosity **significantly**

